

Dark Matter Computations

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Contents

photons

$$\Omega_{\gamma} = 0.005 \%$$

baryons

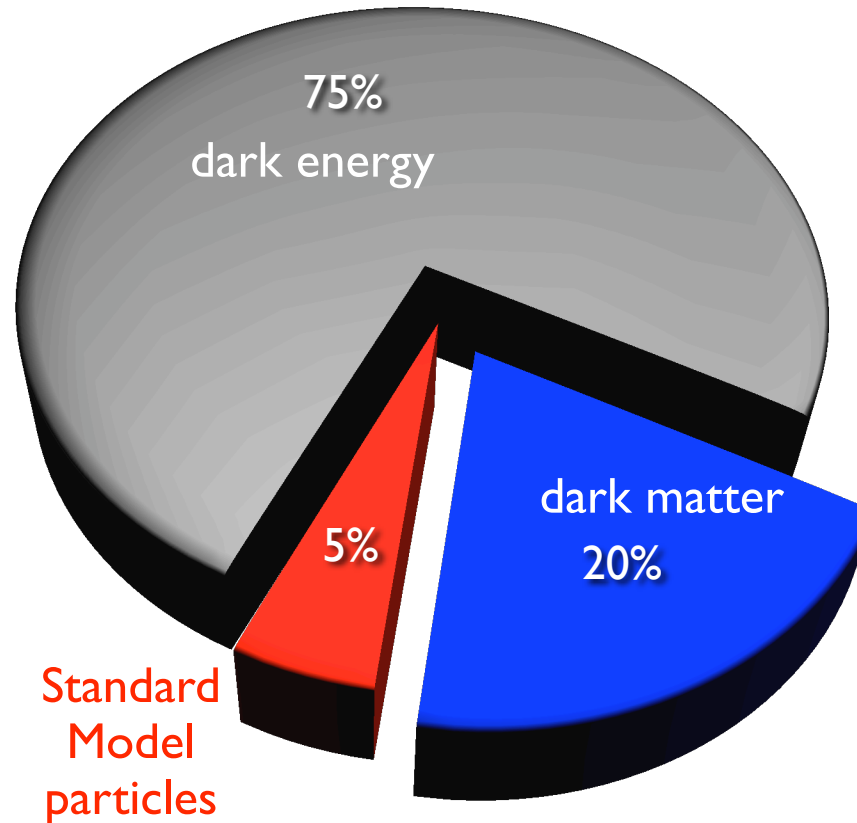
$$\Omega_{\text{B}} = 4 \%$$

? baryon asymmetry ?

neutrinos

$$0.1 \% \leq \Omega_{\nu} \leq 1.5 \%$$

? neutrino mass ?



dark energy

$$\Omega_{\text{DE}} = 75 \%$$

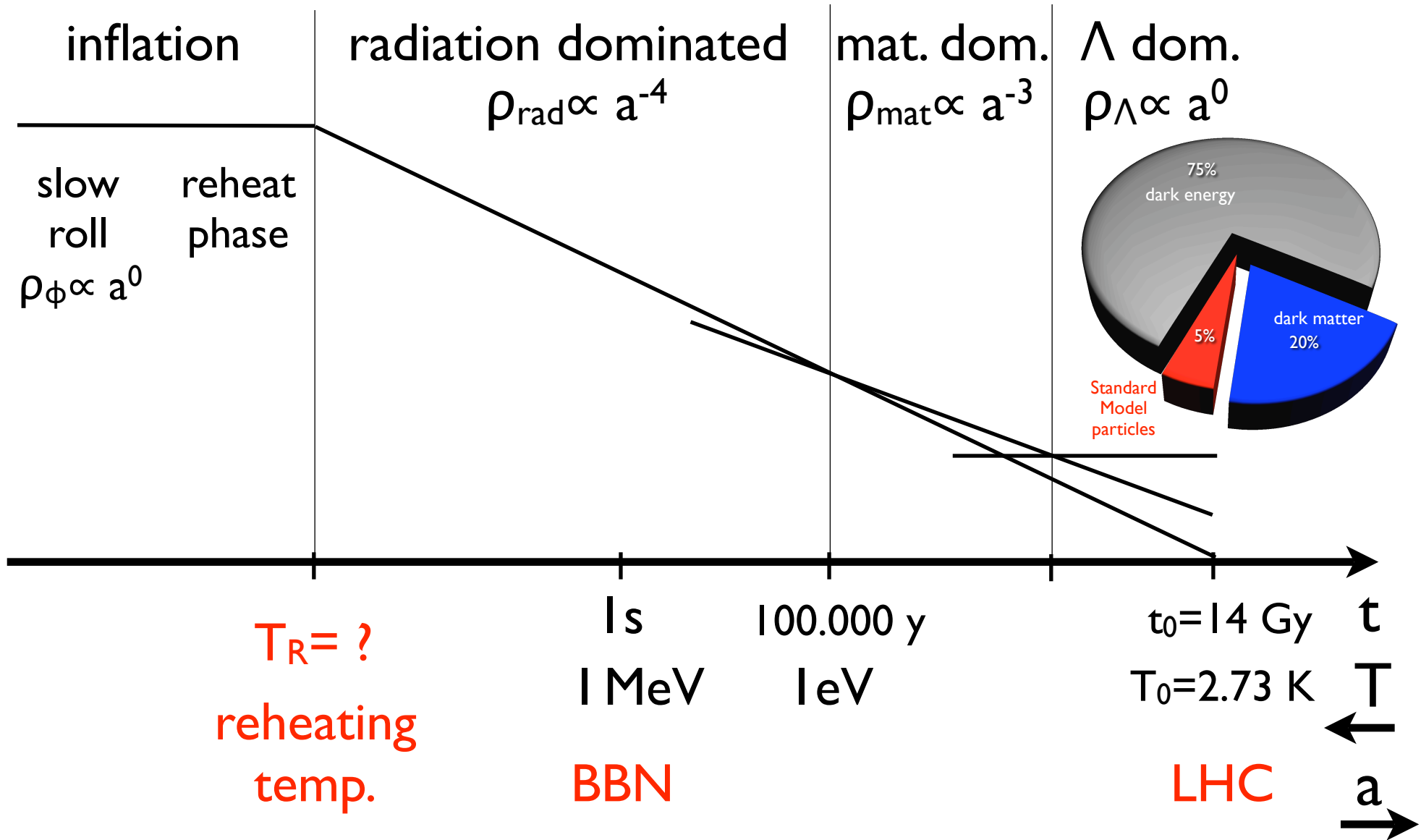
? vacuum energy ?

dark matter

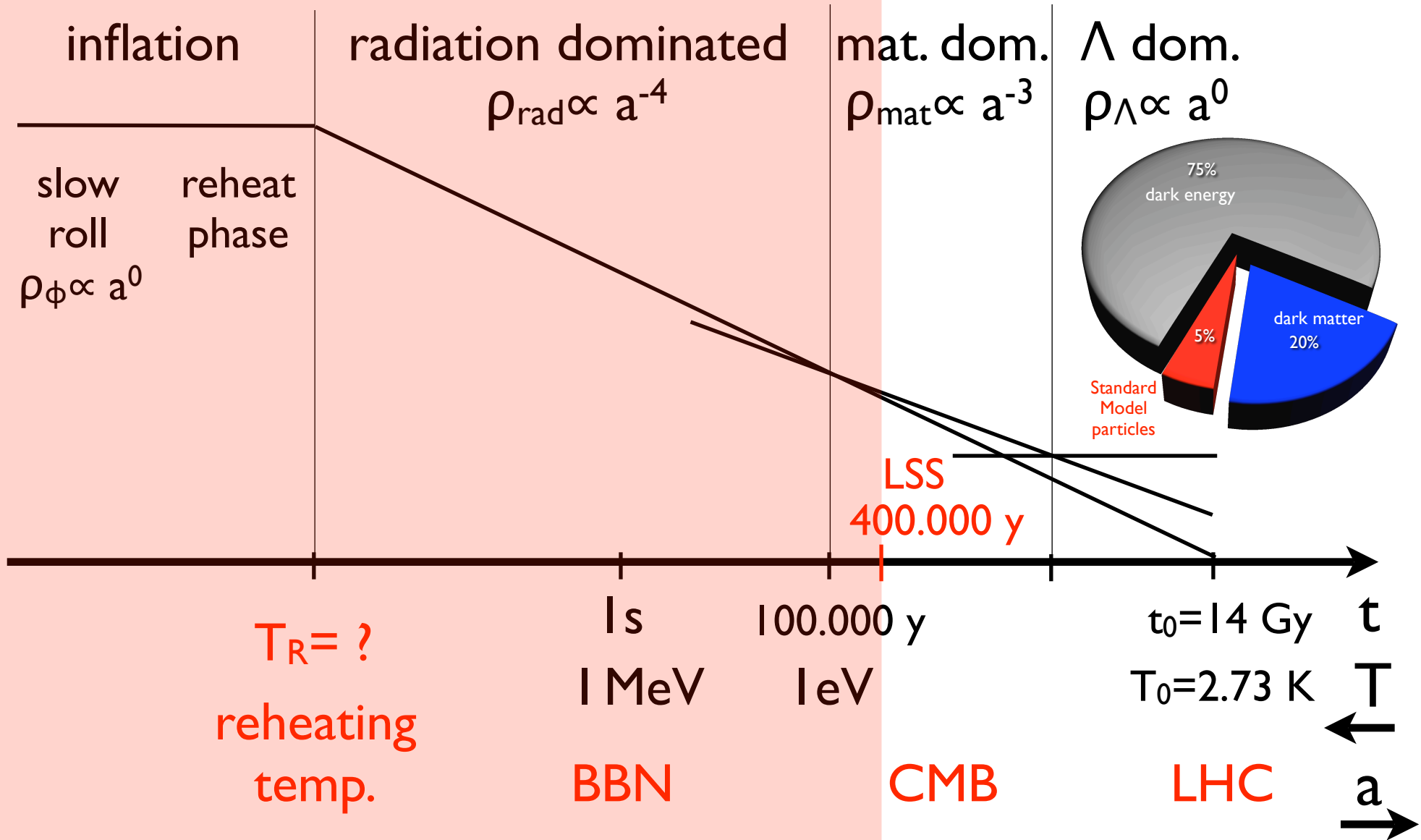
$$\Omega_{\text{DM}} = 20 \%$$

? identity ?

Standard Thermal History of the Universe



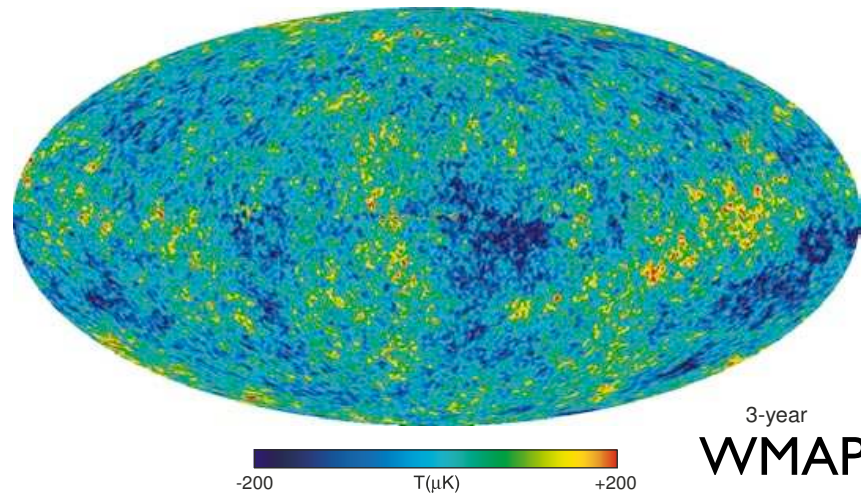
Standard Thermal History of the Universe



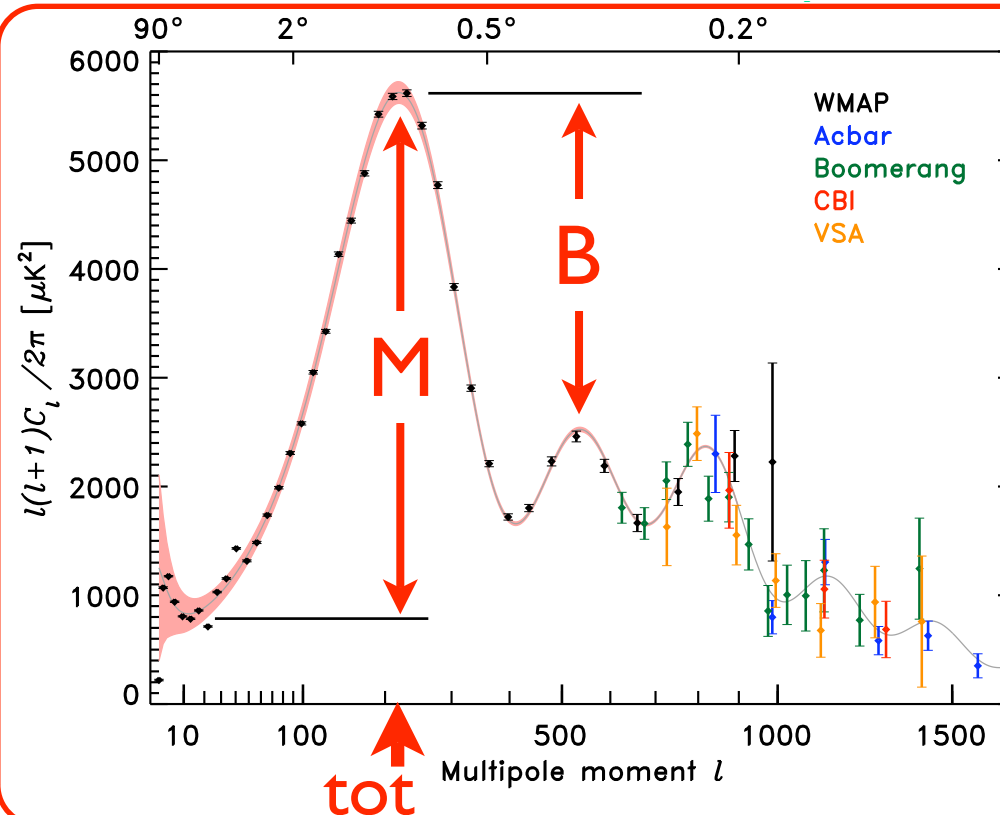
Photons

$\Omega_{\gamma} = 0.005 \%$
photon density

Cosmic
Microwave
Background
Radiation



**Looking
forward
to see
data from
Planck!!!**

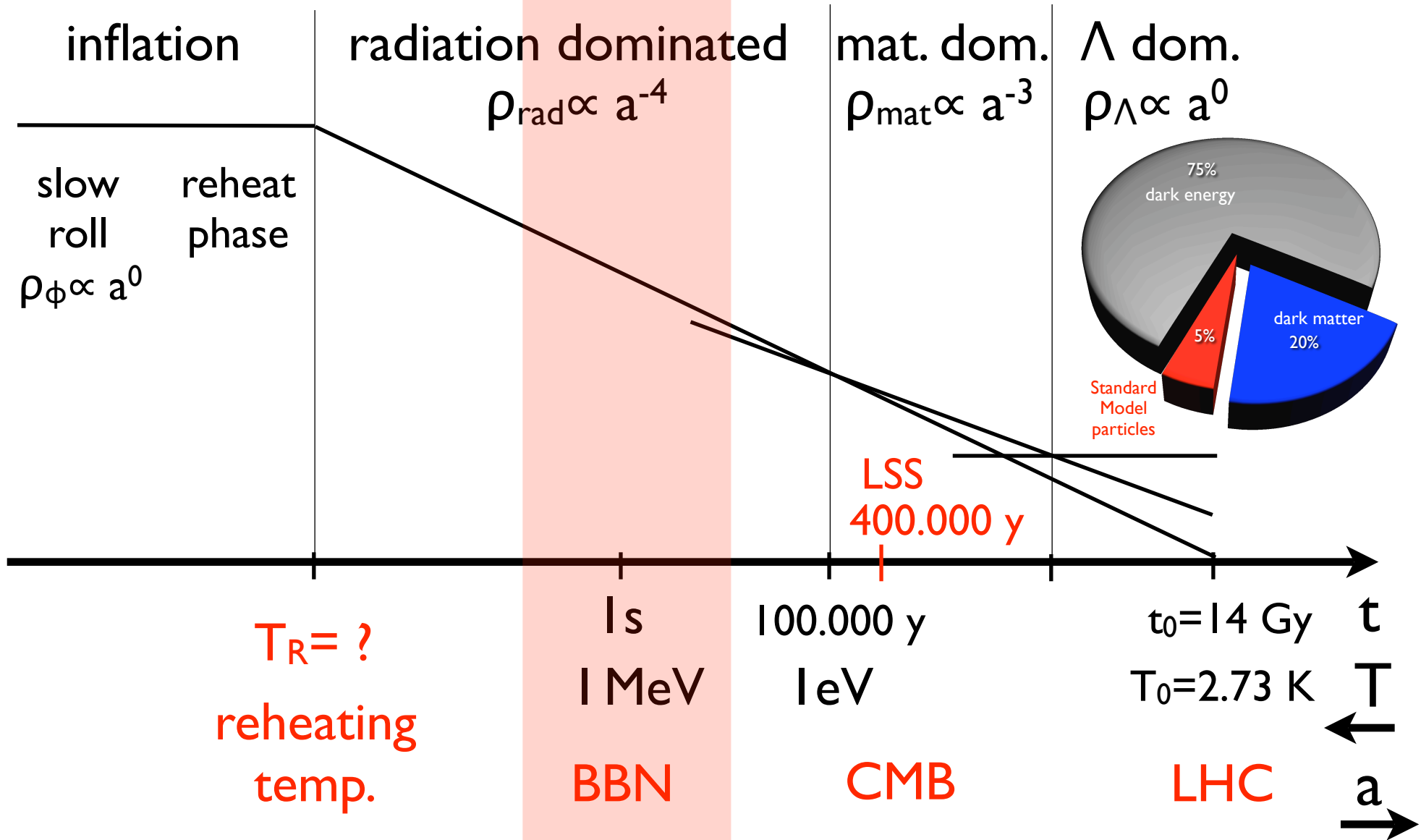


$\Omega_{\text{tot}} = 100 \%$
critical density

$\Omega_{\text{M}} = 24 \%$
matter density

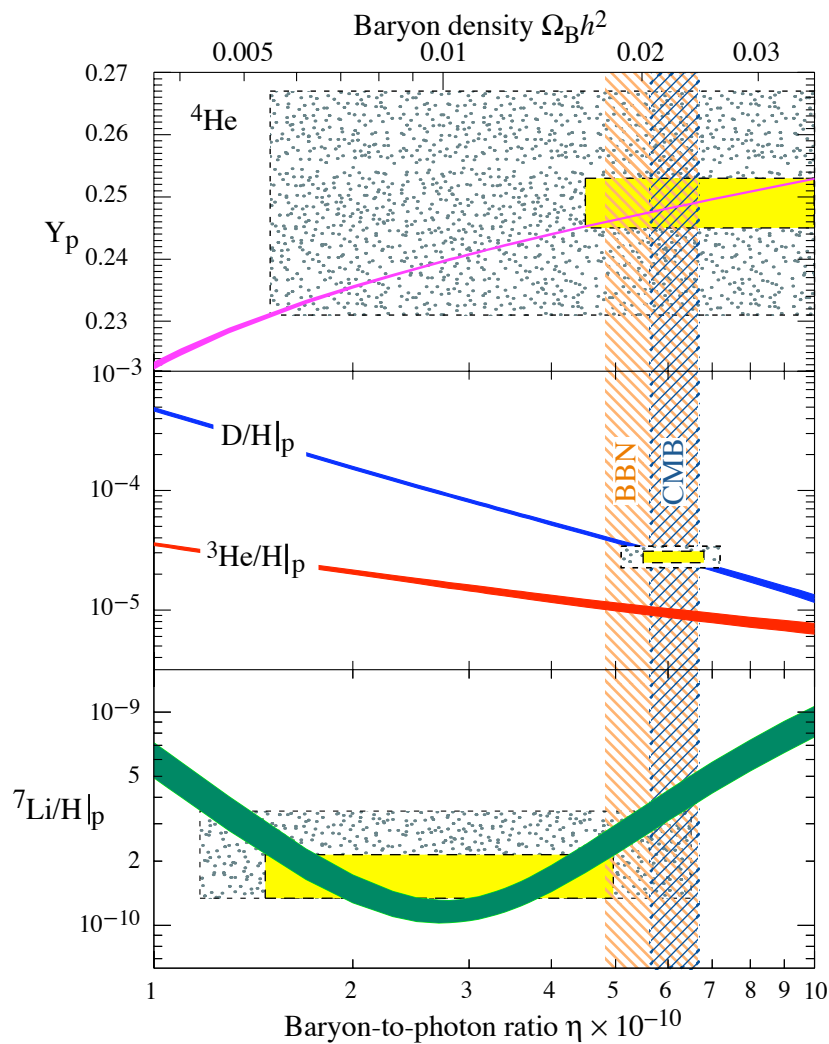
$\Omega_{\text{B}} = 4 \%$
baryon density

Standard Thermal History of the Universe

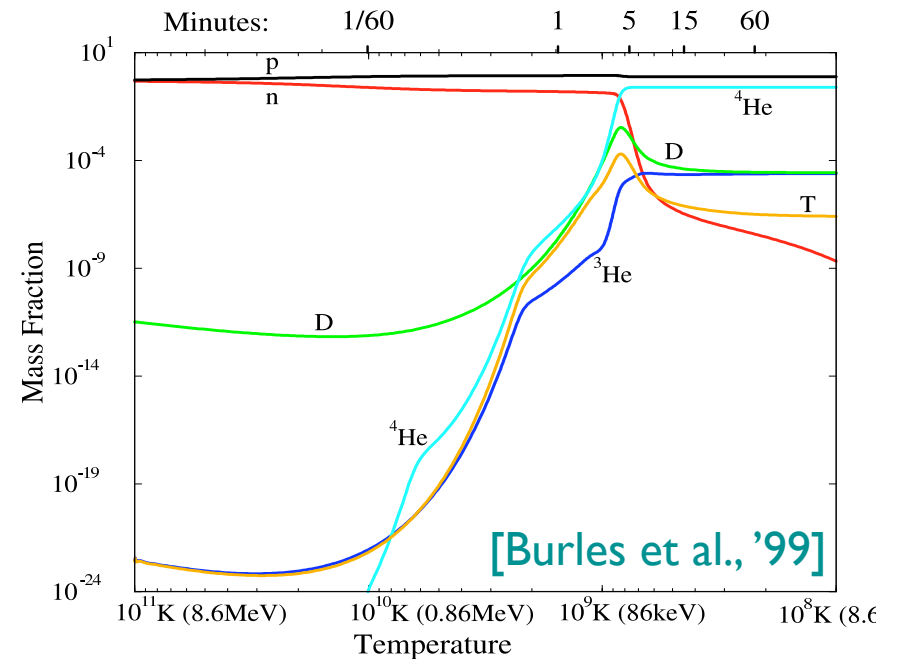


Baryons

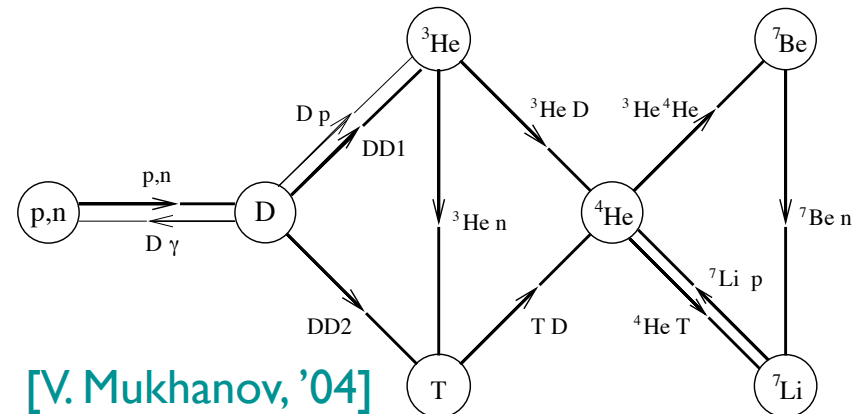
Big Bang Nucleosynthesis



[Particle Data Book 2006]



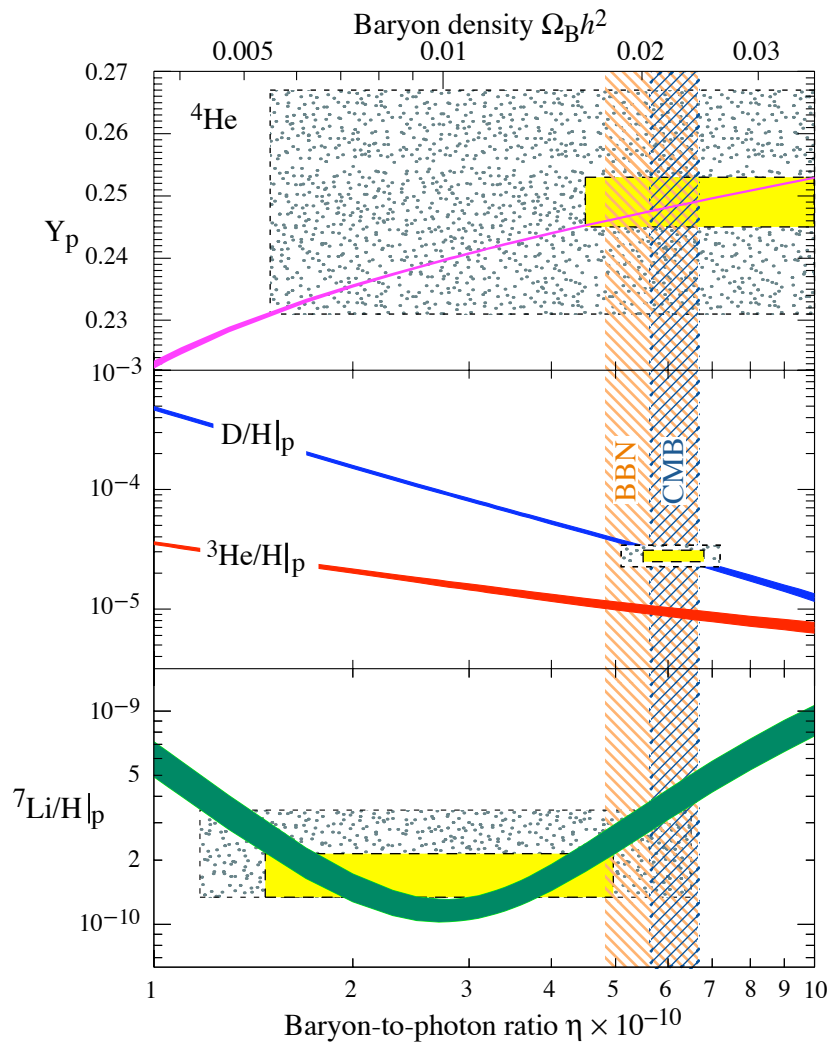
[Burles et al., '99]



[V. Mukhanov, '04]

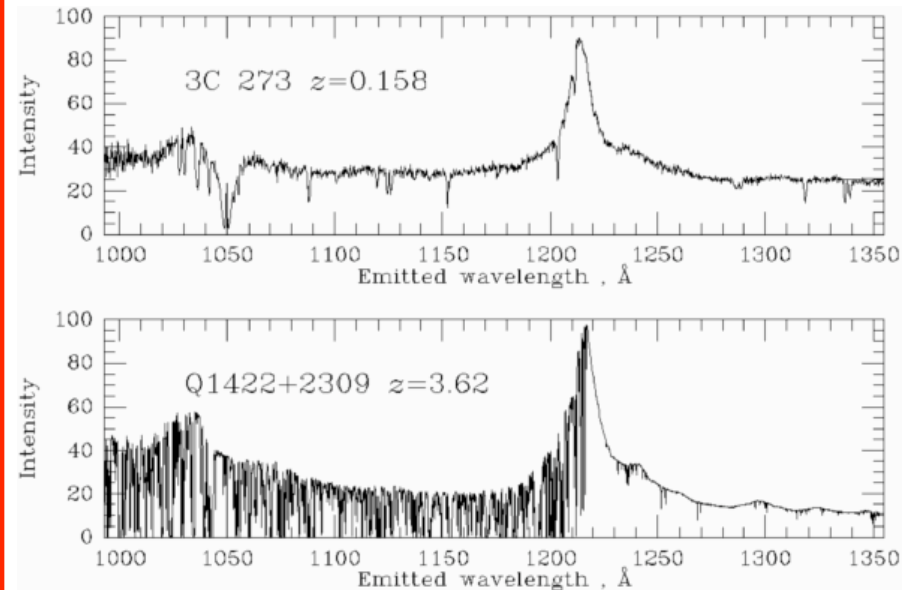
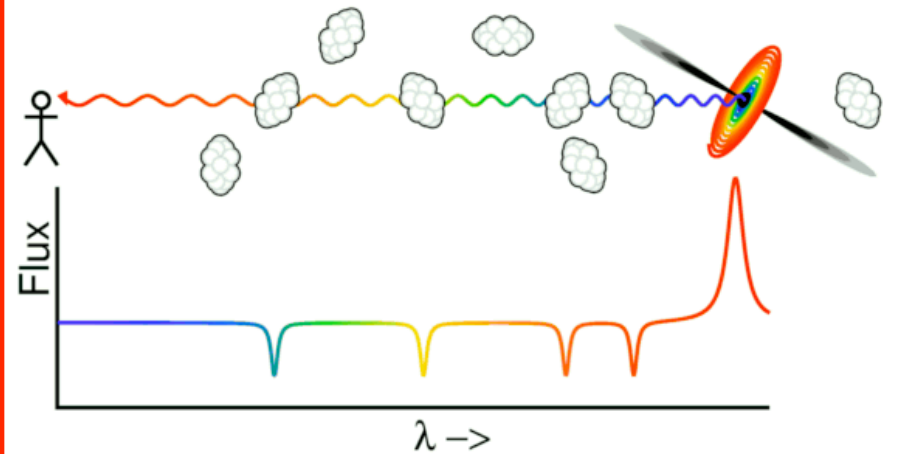
Baryons

Big Bang Nucleosynthesis

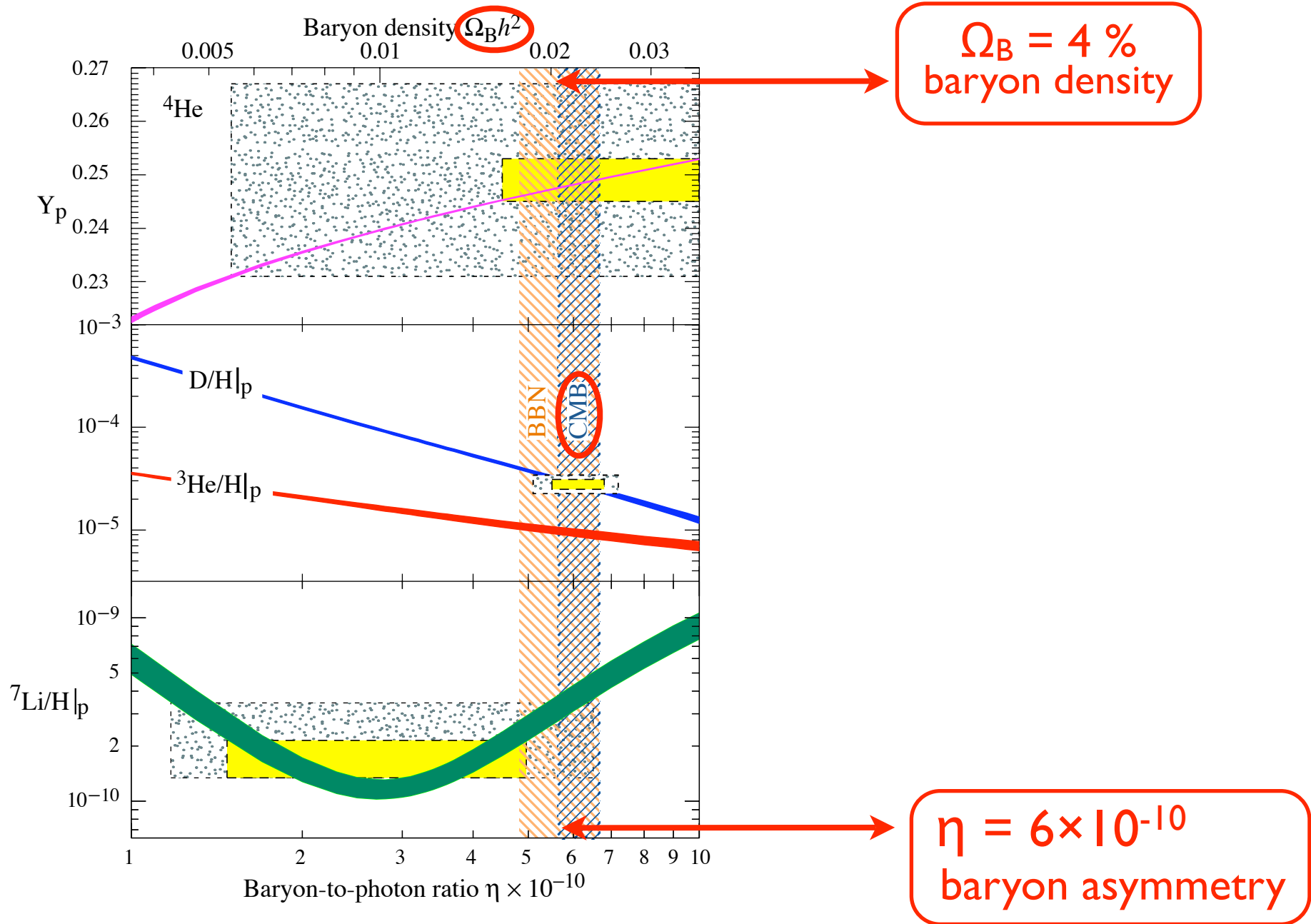


[Particle Data Book 2006]

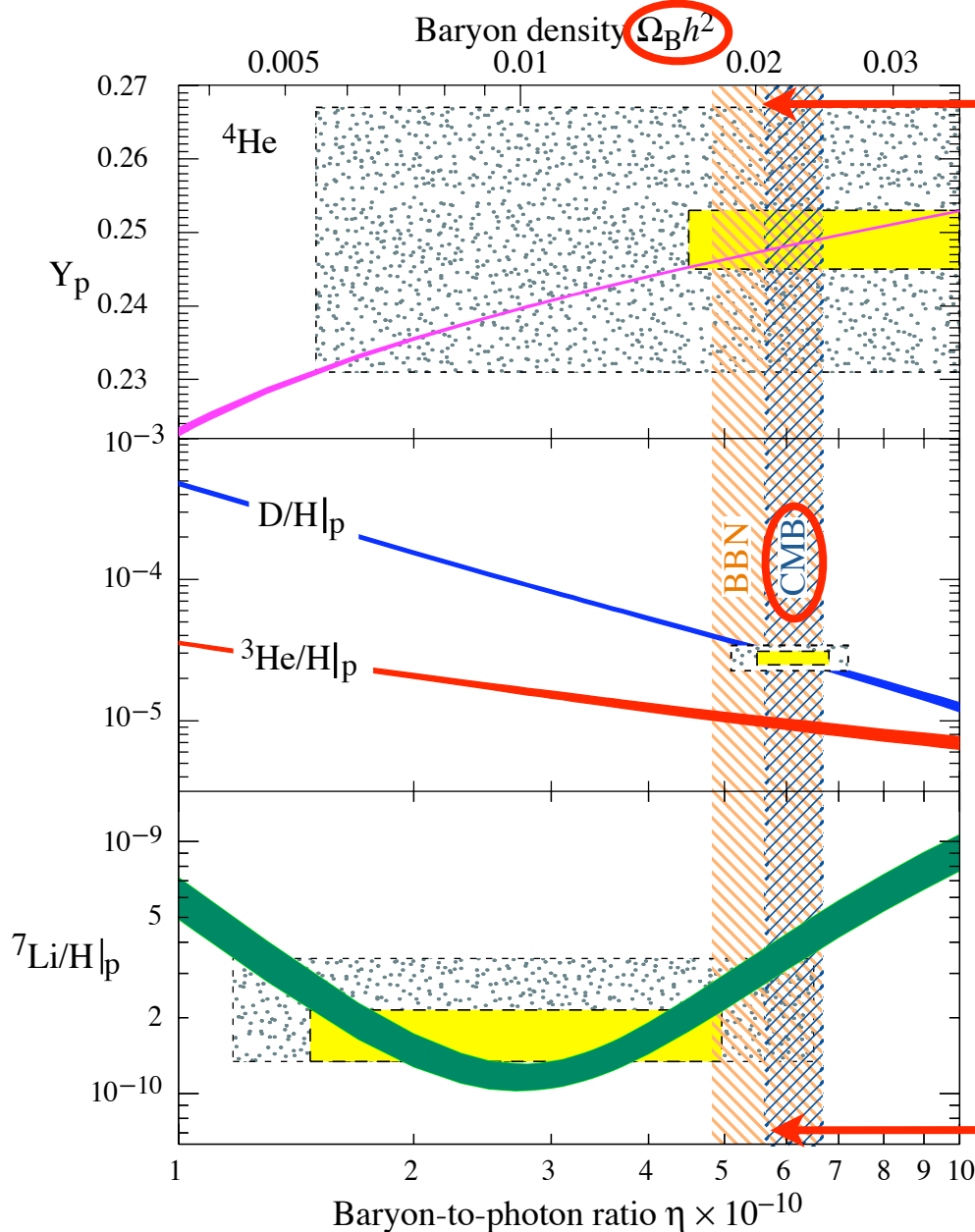
Lyman α Forrest



Baryons



Baryons



$\Omega_B = 4 \%$
baryon density

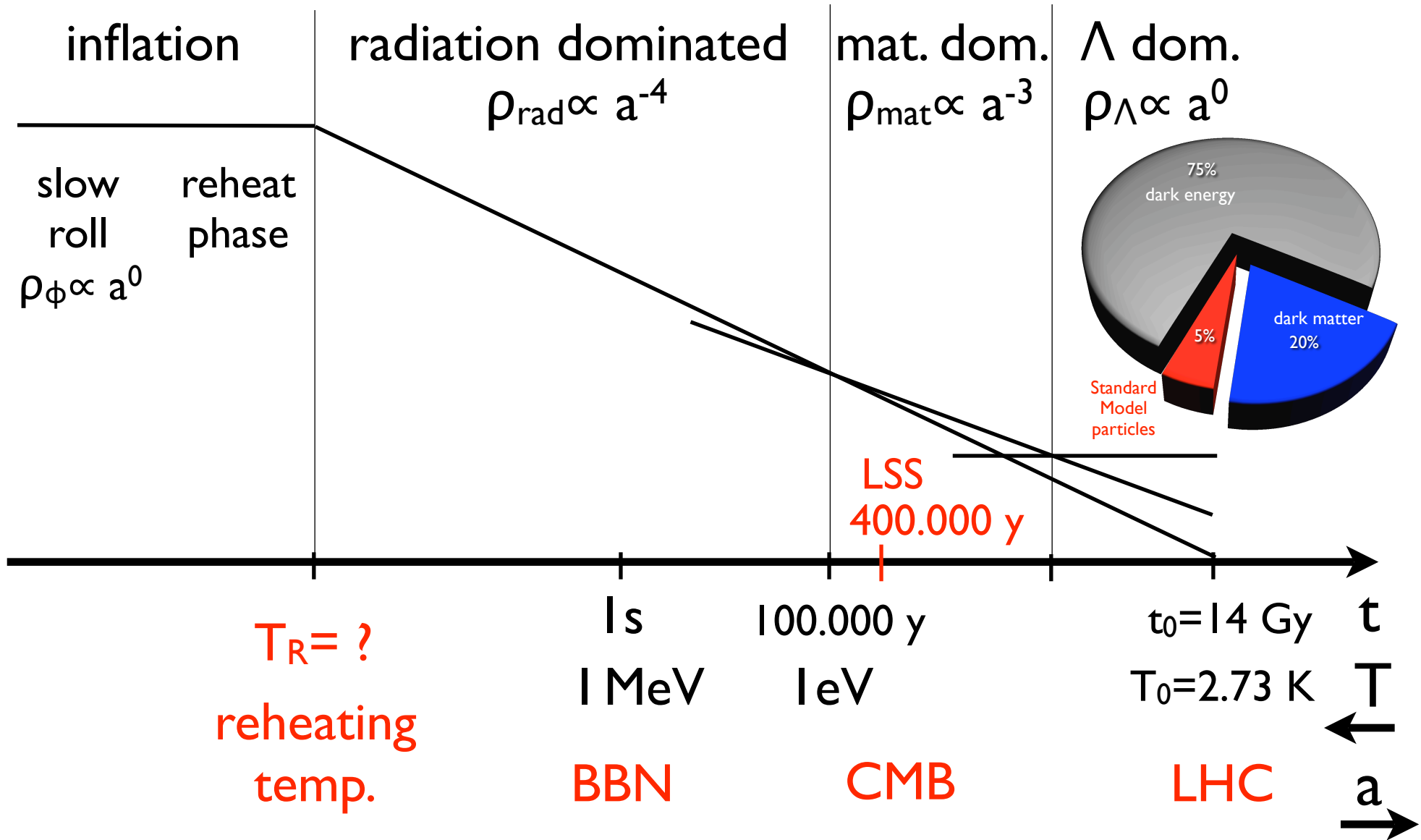
**What is the origin
of the
baryon asymmetry ?**

Electroweak Baryogenesis
 $m_H < 120 \text{ GeV}$
[Talks by Carlos Wagner]

Thermal Leptogenesis
 $T_R > 10^9 \text{ GeV}$

$\eta = 6 \times 10^{-10}$
baryon asymmetry

Standard Thermal History of the Universe

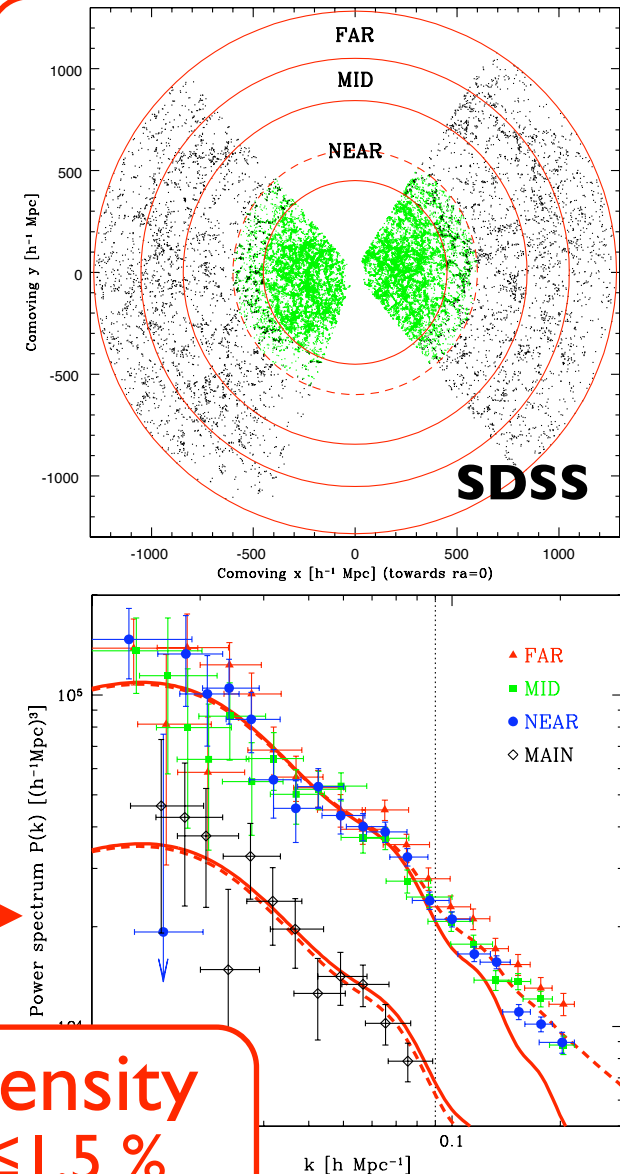


Neutrinos

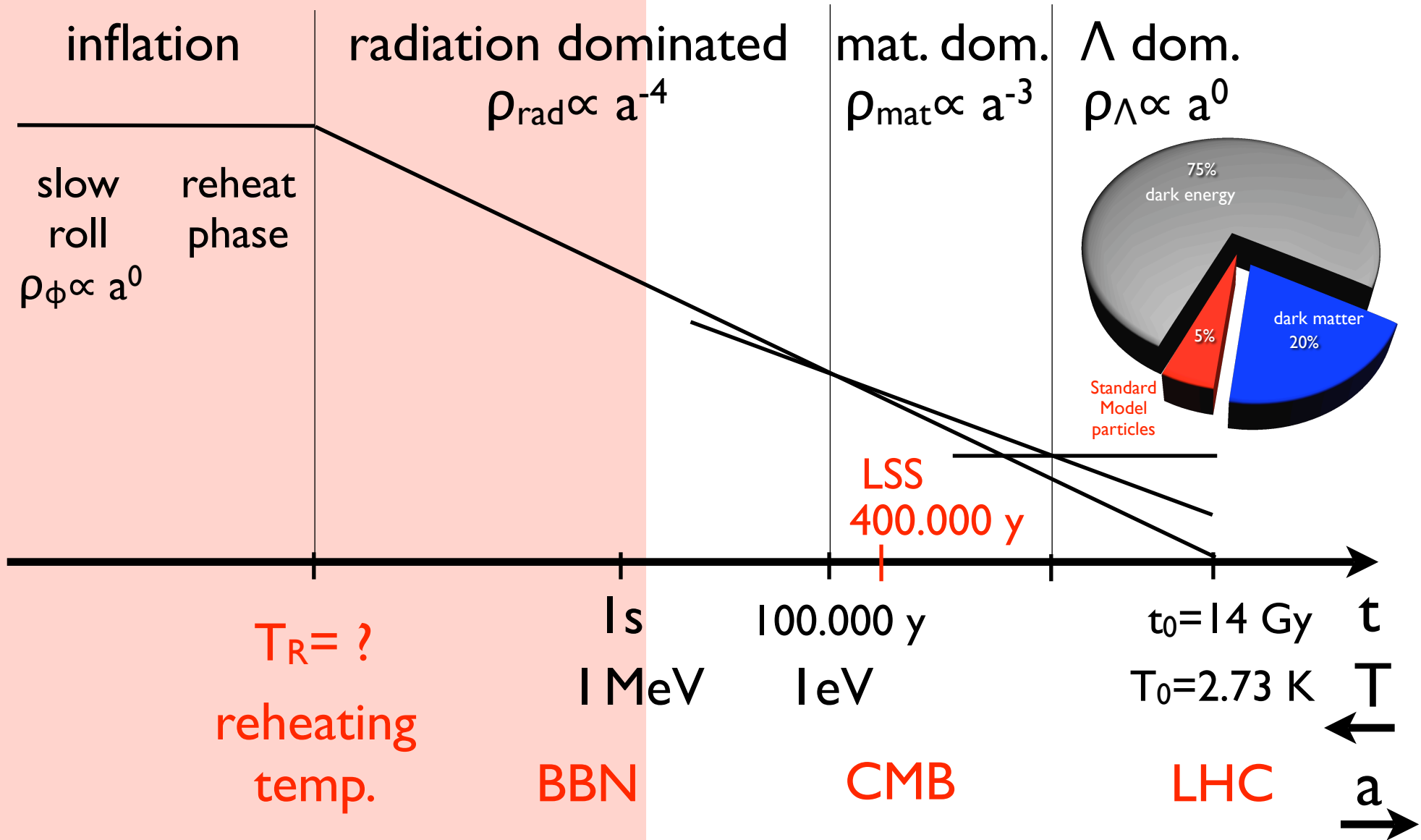
- weakly interacting
- hot thermal relics
- neutrino oscillations
- hot dark matter

$$\sum_i m_{\nu_i} \lesssim \mathcal{O}(1 \text{ eV})$$

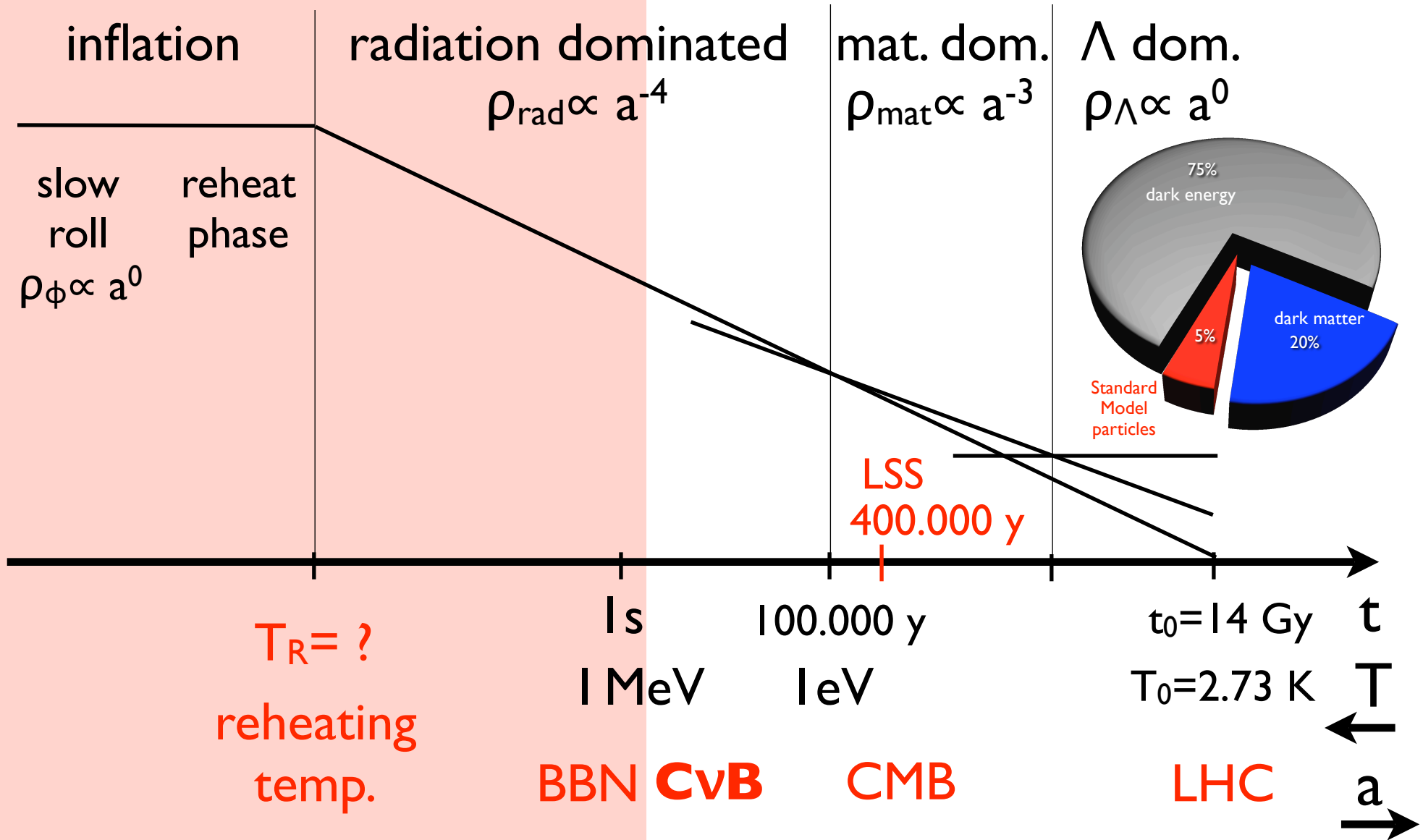
neutrino density
 $0.1 \% \leq \Omega_\nu \leq 1.5 \%$



Standard Thermal History of the Universe



Standard Thermal History of the Universe



Dark Matter

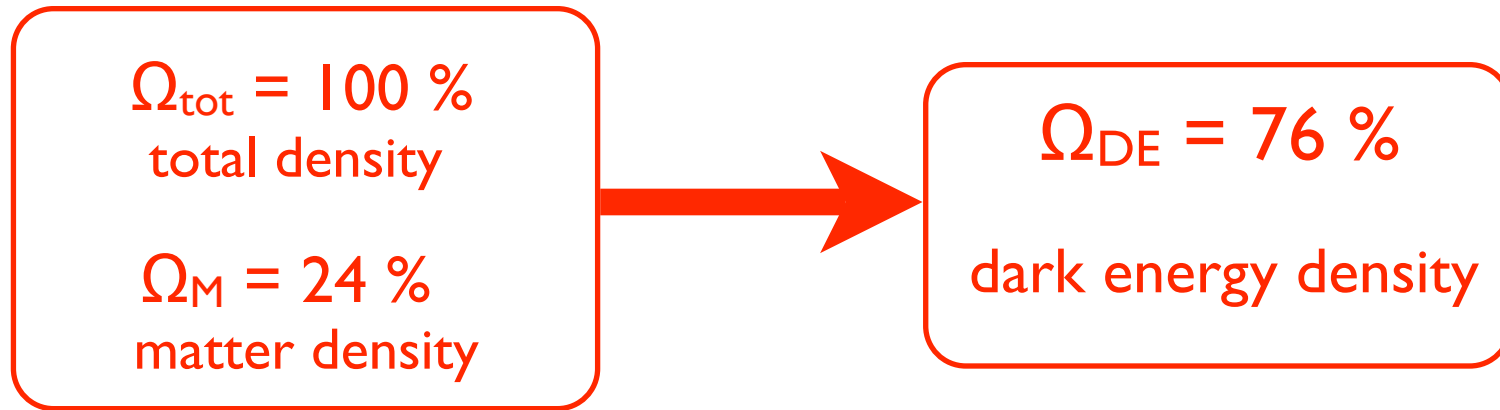
$\Omega_M = 24 \%$
matter density

$\Omega_B = 4 \%$
baryon density



$\Omega_{DM} = 20 \%$
dark matter density

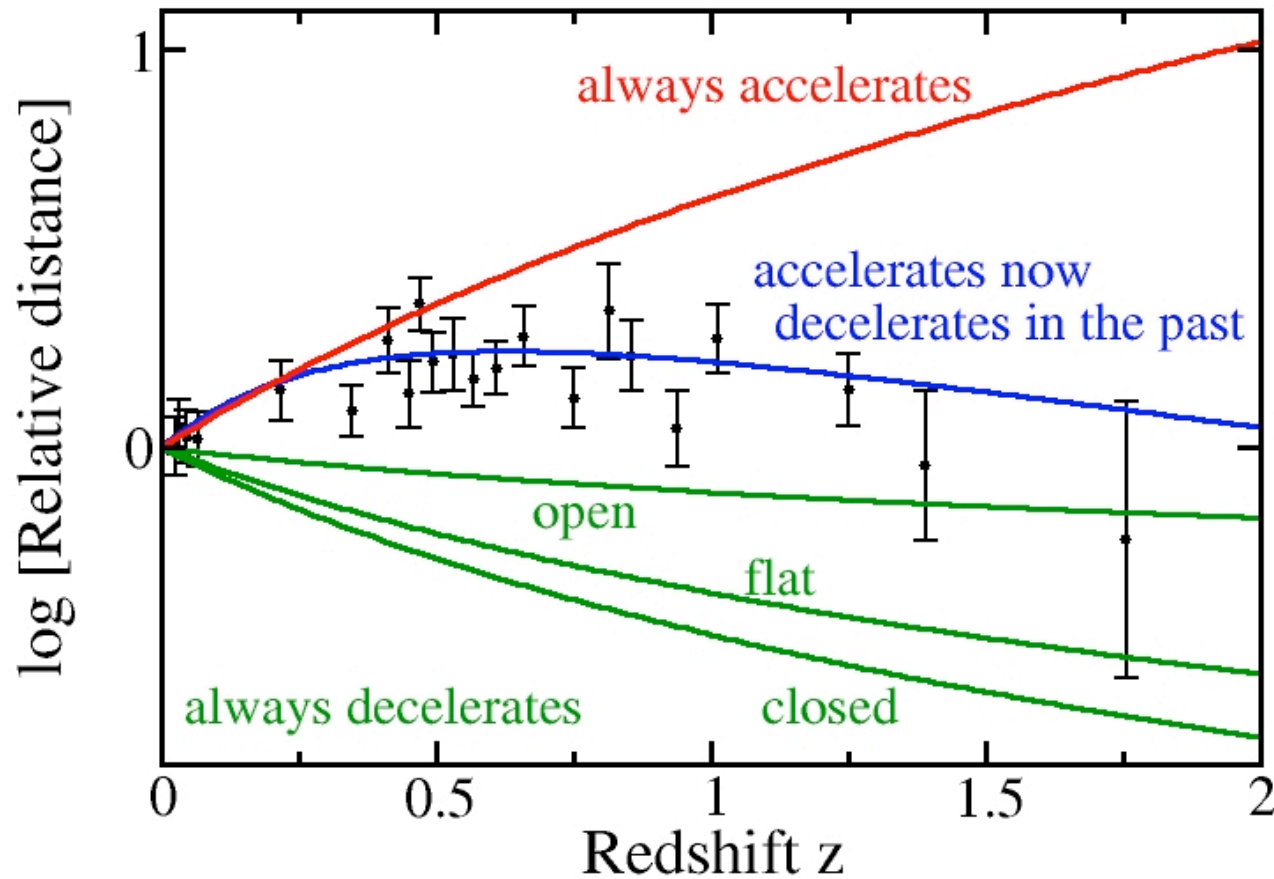
Dark Energy



Cosmological Constant Problem

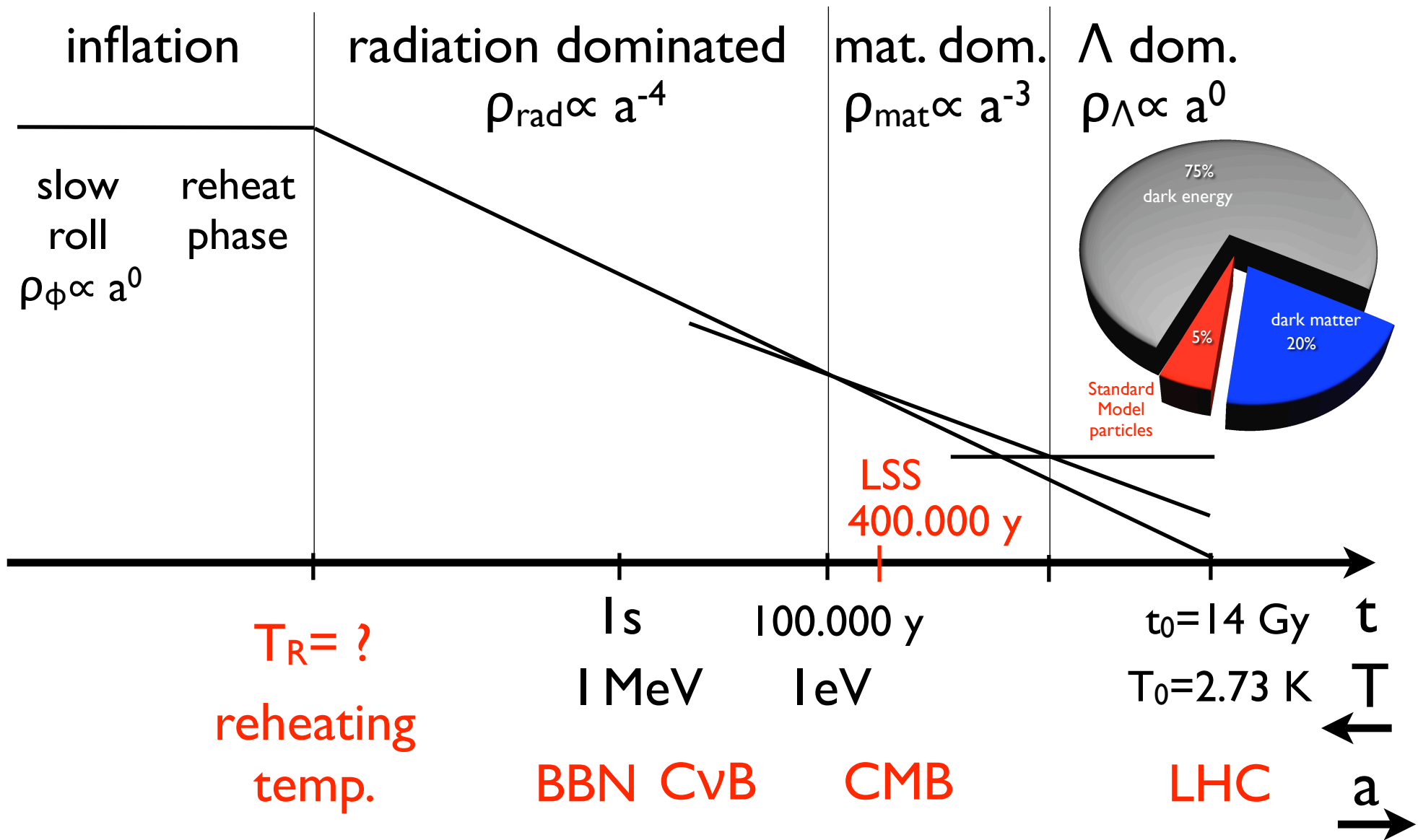
Dark Energy

Supernovae Ia

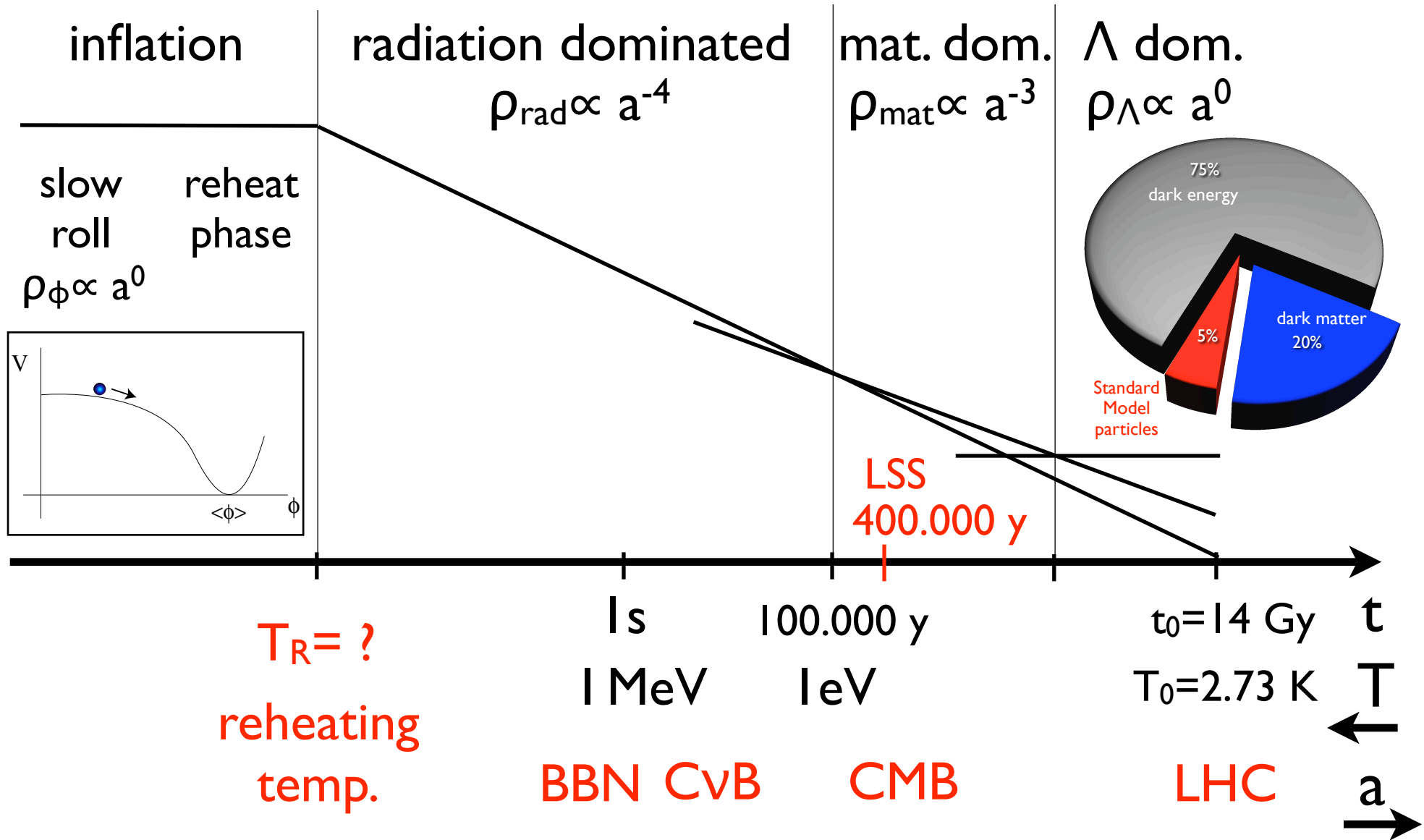


[from Huterer's Talk at PONTAvignon 2008]

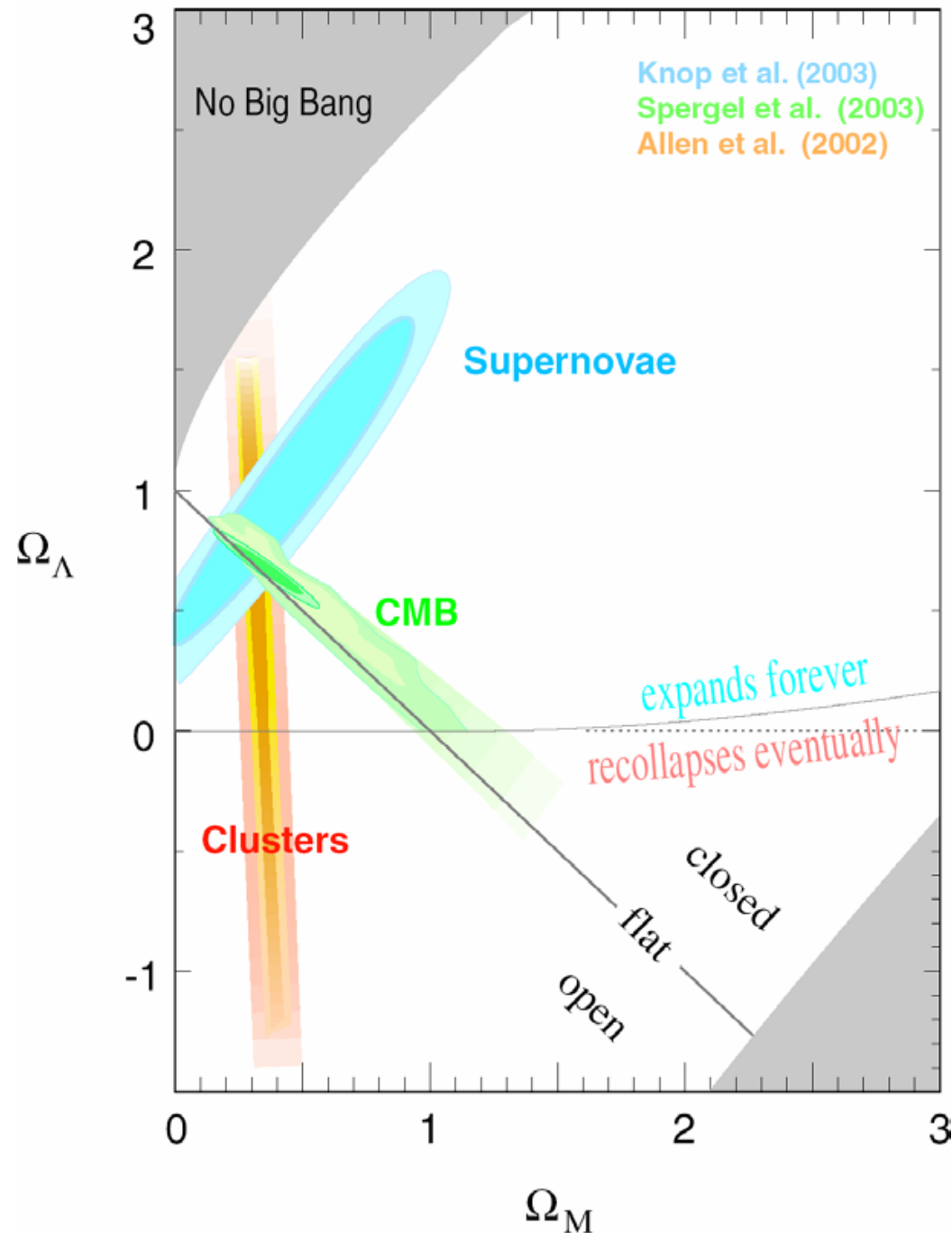
Standard Thermal History of the Universe



Standard Thermal History of the Universe



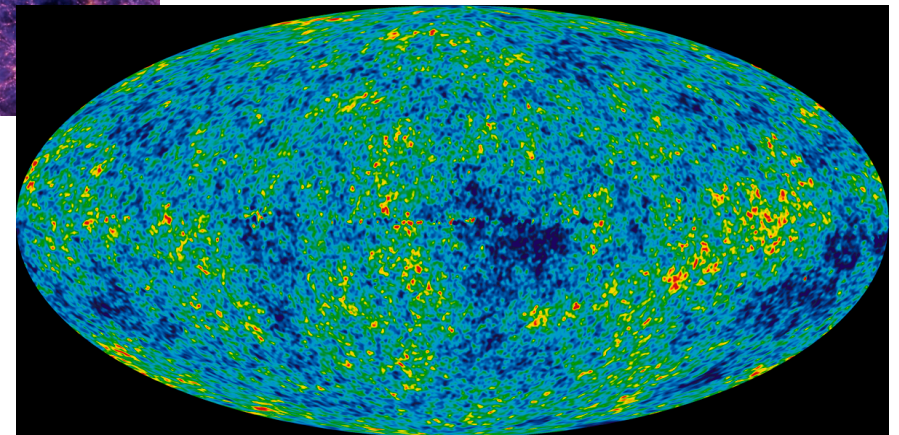
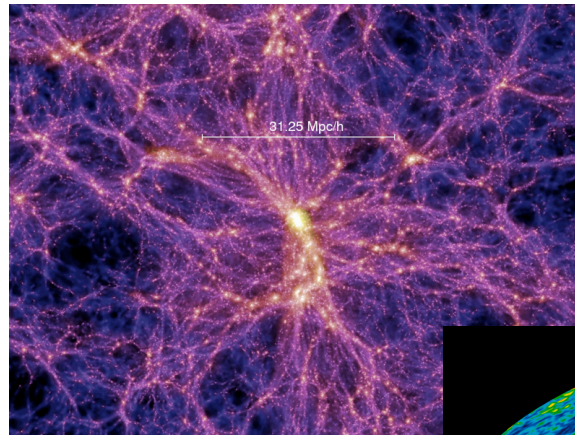
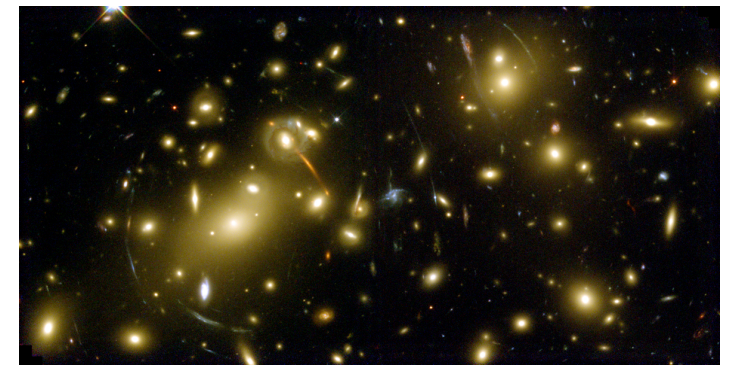
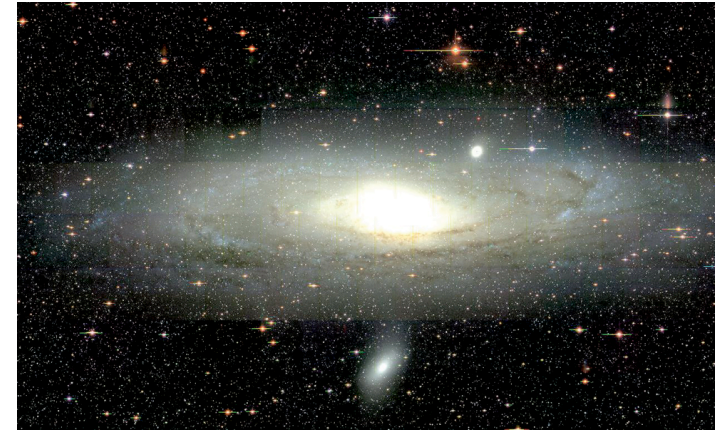
The Cosmic Concordance Model



**What is
the (particle ?)
identity
of dark matter???**

Properties of Dark Matter

- stable or lifetime well above the age of our Universe
- electrically neutral
- clusters →
- “cold”
- dissipationless
- color neutral



The Standard Model

GAUGE	Gauge bosons	$(\text{SU}(3)_c, \text{SU}(2)_L)_Y$
B-boson	$A_\mu^{(1)} = B_\mu$	$(\mathbf{1}, \mathbf{1})_0$
W-bosons	$A_\mu^{(2) a} = W_\mu^a$	$(\mathbf{1}, \mathbf{3})_0$
gluon	$A_\mu^{(3) a} = G_\mu^a$	$(\mathbf{8}, \mathbf{1})_0$
MATTER	Fermions	$(\text{SU}(3)_c, \text{SU}(2)_L)_Y$
leptons $I = 1, 2, 3$	$L^I = \begin{pmatrix} \nu_L^I \\ e_L^{-I} \end{pmatrix}$	$(\mathbf{1}, \mathbf{2})_{-1}$
	$E^{c I} = e_R^{-c I}$	$(\mathbf{1}, \mathbf{1})_{+2}$
quarks $I = 1, 2, 3$ ($\times 3$ colors)	$Q^I = \begin{pmatrix} u_L^I \\ d_L^I \end{pmatrix}$	$(\mathbf{3}, \mathbf{2})_{+\frac{1}{3}}$
	$U^{c I} = u_R^{c I}$	$(\bar{\mathbf{3}}, \mathbf{1})_{-\frac{4}{3}}$
	$D^{c I} = d_R^{c I}$	$(\bar{\mathbf{3}}, \mathbf{1})_{+\frac{2}{3}}$
HIGGS	Higgs Boson	$(\text{SU}(3)_c, \text{SU}(2)_L)_Y$
Higgs	$\phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix}$	$(\mathbf{1}, \mathbf{2})_{+1}$

Dark Matter



Physics beyond the Standard Model

Supersymmetry

GAUGE	Gauge bosons	Gauginos	$(SU(3)_c, SU(2)_L)_Y$
B-boson, bino	$A_\mu^{(1)} = B_\mu$	$\lambda^{(1)} = \tilde{B}$	$(1, 1)_0$
W-bosons, winos	$A_\mu^{(2)a} = W_\mu^a$	$\lambda^{(2)a} = \tilde{W}^a$	$(1, 3)_0$
gluon, gluino	$A_\mu^{(3)a} = G_\mu^a$	$\lambda^{(3)a} = \tilde{g}^a$	$(8, 1)_0$
MATTER	Sfermions	Fermions	$(SU(3)_c, SU(2)_L)_Y$
sleptons, leptons $I = 1, 2, 3$	$\tilde{L}^I = \begin{pmatrix} \tilde{\nu}_L^I \\ \tilde{e}_L^{-I} \end{pmatrix}$ $\tilde{E}^{*I} = \tilde{e}_R^{-*I}$	$L^I = \begin{pmatrix} \nu_L^I \\ e_L^{-I} \end{pmatrix}$ $E^{cI} = e_R^{-cI}$	$(1, 2)_{-1}$ $(1, 1)_{+2}$
squarks, quarks $I = 1, 2, 3$ ($\times 3$ colors)	$\tilde{Q}^I = \begin{pmatrix} \tilde{u}_L^I \\ \tilde{d}_L^I \end{pmatrix}$ $\tilde{U}^{*I} = \tilde{u}_R^{*I}$ $\tilde{D}^{*I} = \tilde{d}_R^{*I}$	$Q^I = \begin{pmatrix} u_L^I \\ d_L^I \end{pmatrix}$ $U^{cI} = u_R^{cI}$ $D^{cI} = d_R^{cI}$	$(3, 2)_{+\frac{1}{3}}$ $(\bar{3}, 1)_{-\frac{4}{3}}$ $(\bar{3}, 1)_{+\frac{2}{3}}$
Higgs, higgsinos	$H_d = \begin{pmatrix} H_d^0 \\ H_d^- \end{pmatrix}$ $H_u = \begin{pmatrix} H_u^+ \\ H_u^0 \end{pmatrix}$	$\tilde{H}_d = \begin{pmatrix} \tilde{H}_d^0 \\ \tilde{H}_d^- \end{pmatrix}$ $\tilde{H}_u = \begin{pmatrix} \tilde{H}_u^+ \\ \tilde{H}_u^0 \end{pmatrix}$	$(1, 2)_{-1}$ $(1, 2)_{+1}$

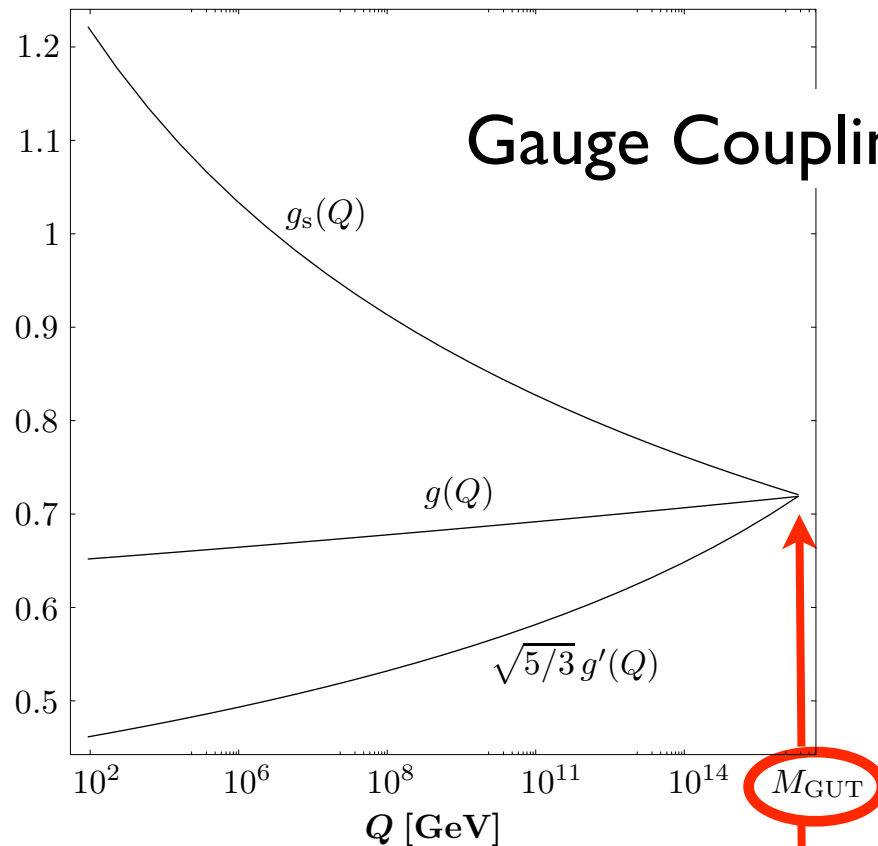
Minimal
Supersymmetric
Extension
of the
Standard Model



Every Particle
of the
Standard Model
has a
Superpartner

Why Supersymmetry?

Extension of Space-Time Symmetry



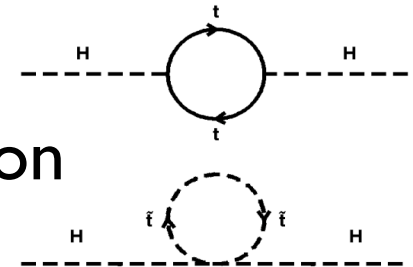
Gauge Coupling Unification

Hierarchy Stabilization

(Super-) Gravity

Consistent String Theory

Dark Matter

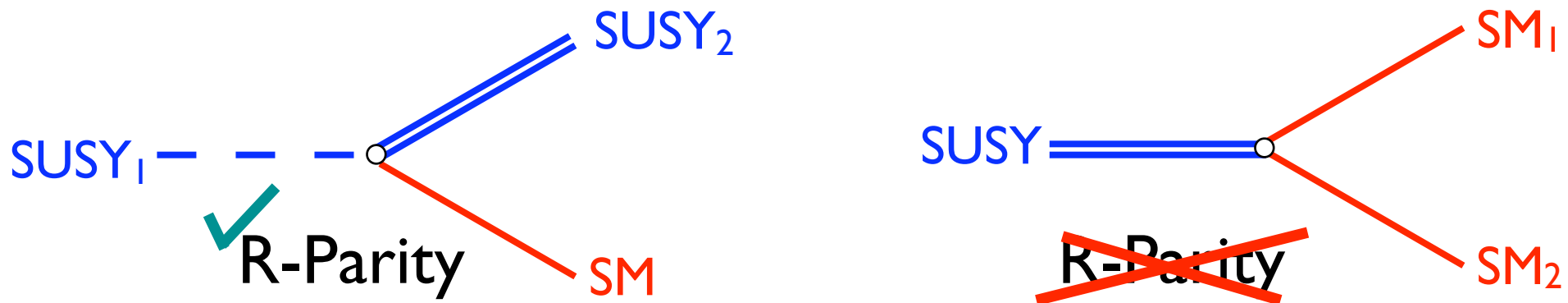


Gauge Coupling Unification at $M_{GUT} \simeq 2 \times 10^{16}$ GeV

Conservation of R-Parity

- superpotential: $W_{\text{MSSM}} \leftarrow W_{\Delta L} + W_{\Delta B}$
- non-observation of L & B violating processes (proton stability, ...)
- postulate conservation of R-Parity \leftarrow multiplicative quantum number

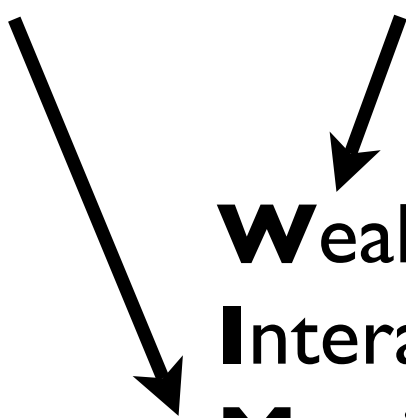
$$P_R = (-1)^{3(B-L)+2S} = \begin{cases} +1 & \text{for SM, } H_u, H_d \\ -1 & \text{for } \tilde{X} \leftarrow \text{superpartners} \end{cases}$$



The lightest supersymmetric particle (LSP) is stable!!!

Supersymmetric Dark Matter Candidates

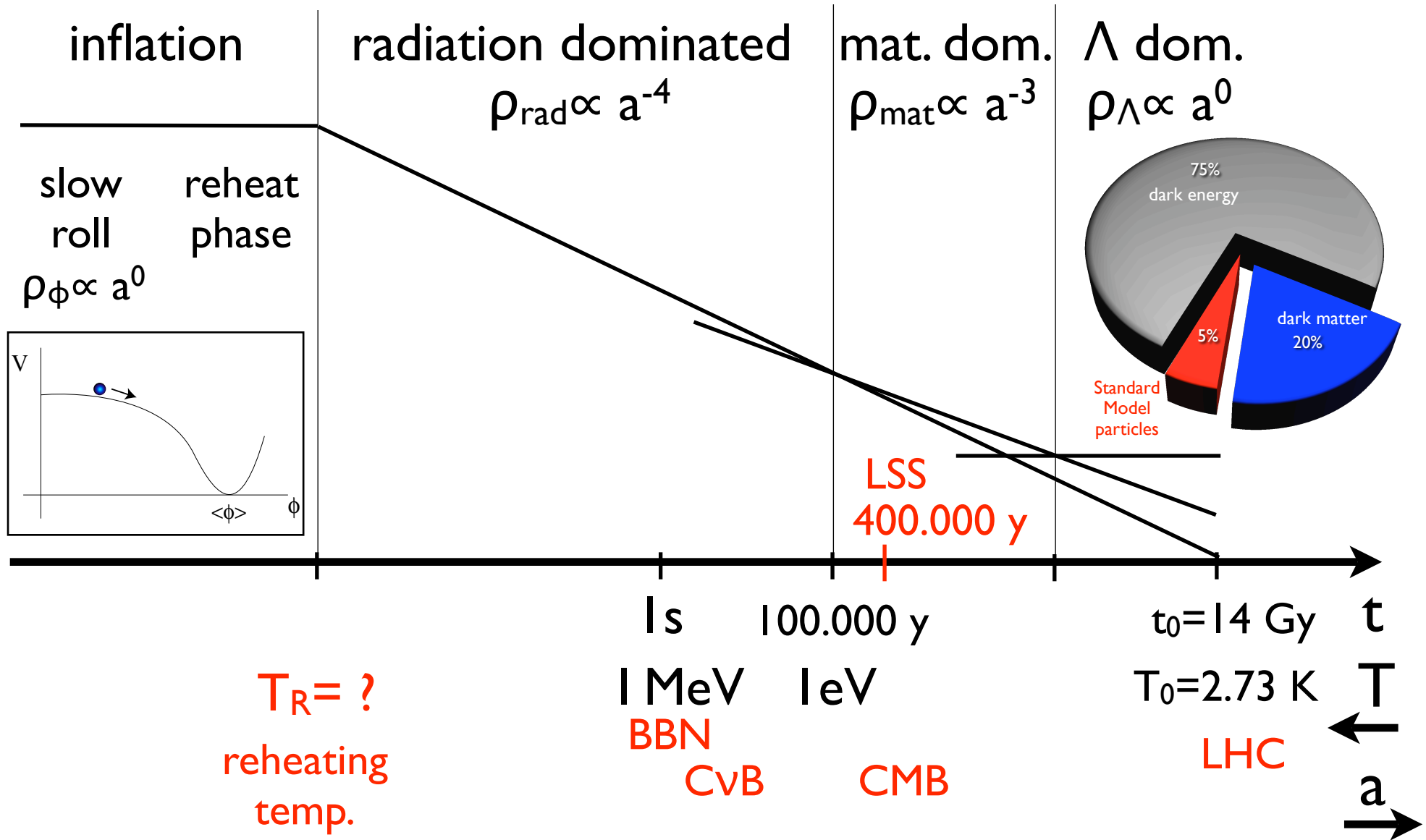
	LSP	ID	spin	mass	interaction
lightest neutralino	$\tilde{\chi}_1^0$	$\tilde{B}, \tilde{W}, \tilde{H}_u^0, \tilde{H}_d^0$	$\frac{1}{2}$	$\mathcal{O}(100 \text{ GeV})$	g, g'
$\in \text{MSSM}$		mixture		$M_1, M_2, \mu, \tan \beta$	weak



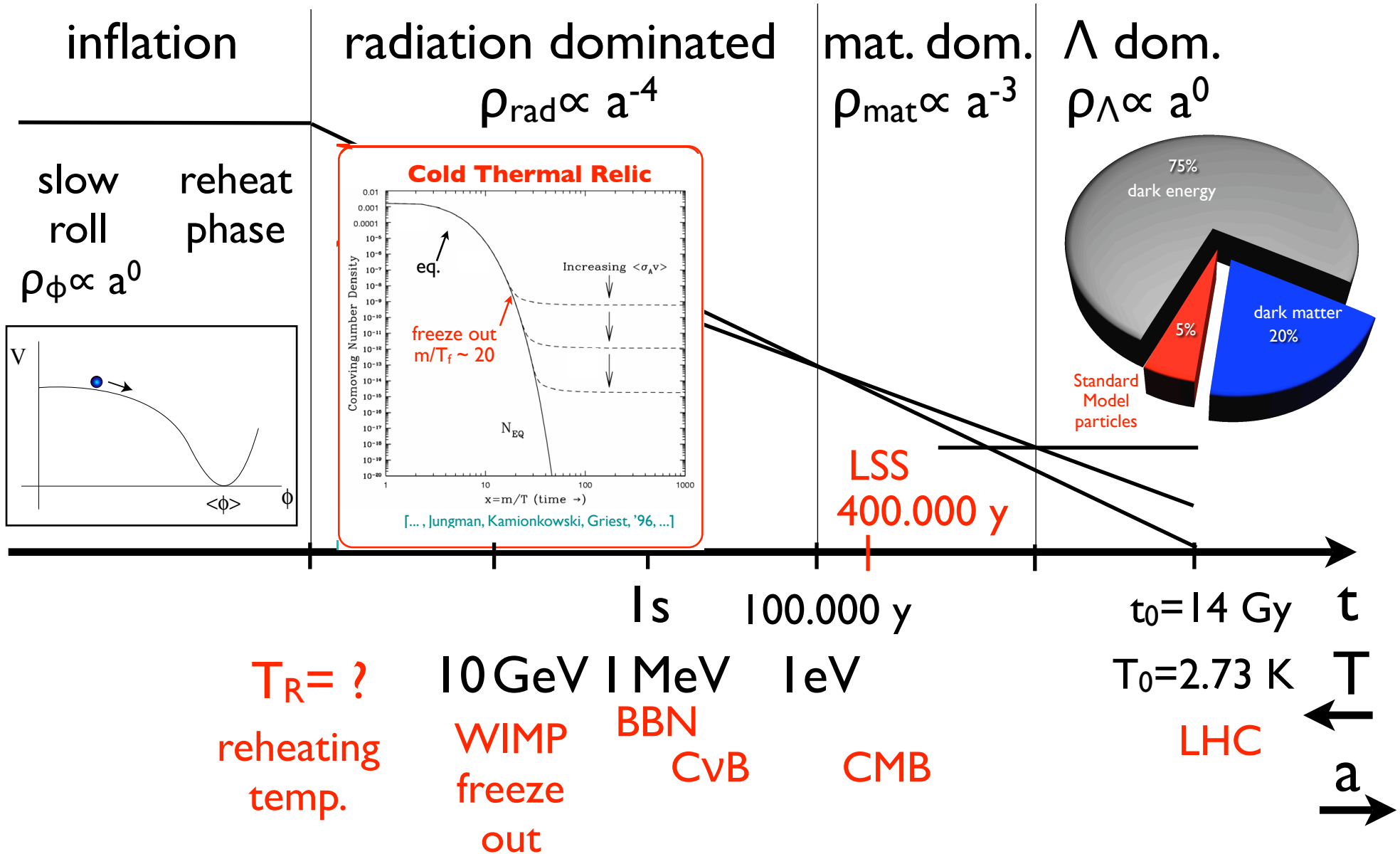
**Weakly
Interacting
Massive
Particle**

["Heavy Neutrino"]

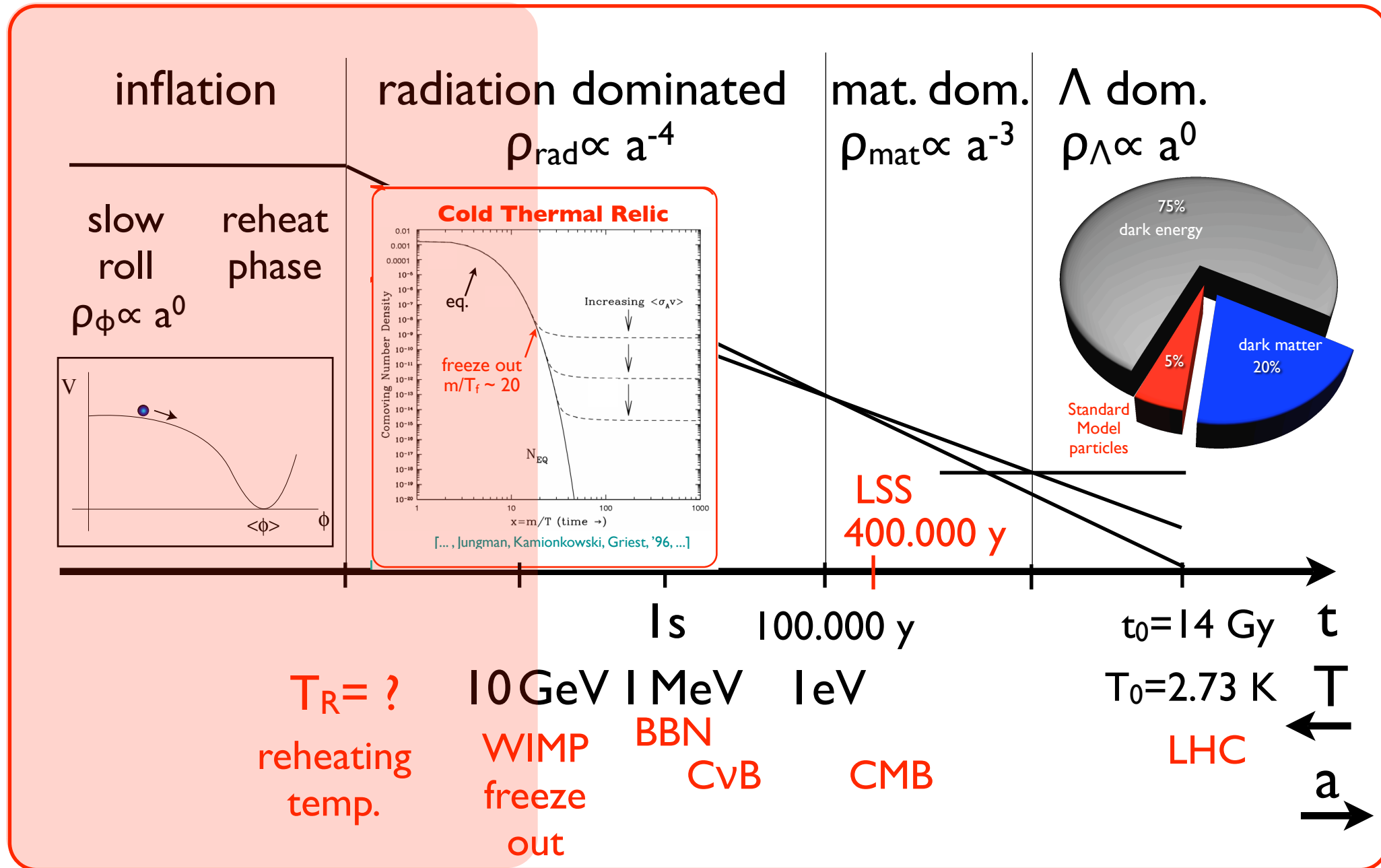
Standard Thermal History of the Universe



Standard Thermal History of the Universe



Standard Thermal History of the Universe



Boltzmann Eqs. for WIMP freeze out

$$\frac{dY}{dT} = \sqrt{\frac{\pi g_*(T)}{45}} M_p \langle \sigma v \rangle (Y(T)^2 - Y_{eq}(T)^2)$$

$$\langle \sigma v \rangle = \frac{\sum_{i,j} g_i g_j \int_{(m_i+m_j)^2} ds \sqrt{s} K_1(\sqrt{s}/T) p_{ij}^2 \sigma_{ij}(s)}{2T \left(\sum_i g_i m_i^2 K_2(m_i/T) \right)^2}$$

$\tilde{\chi}_1^0$ LSP Dark Matter: Production, Constraints

LSP	interaction	production	constraints
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$\tilde{\chi}_1^0$

g, g'

weak

WIMP

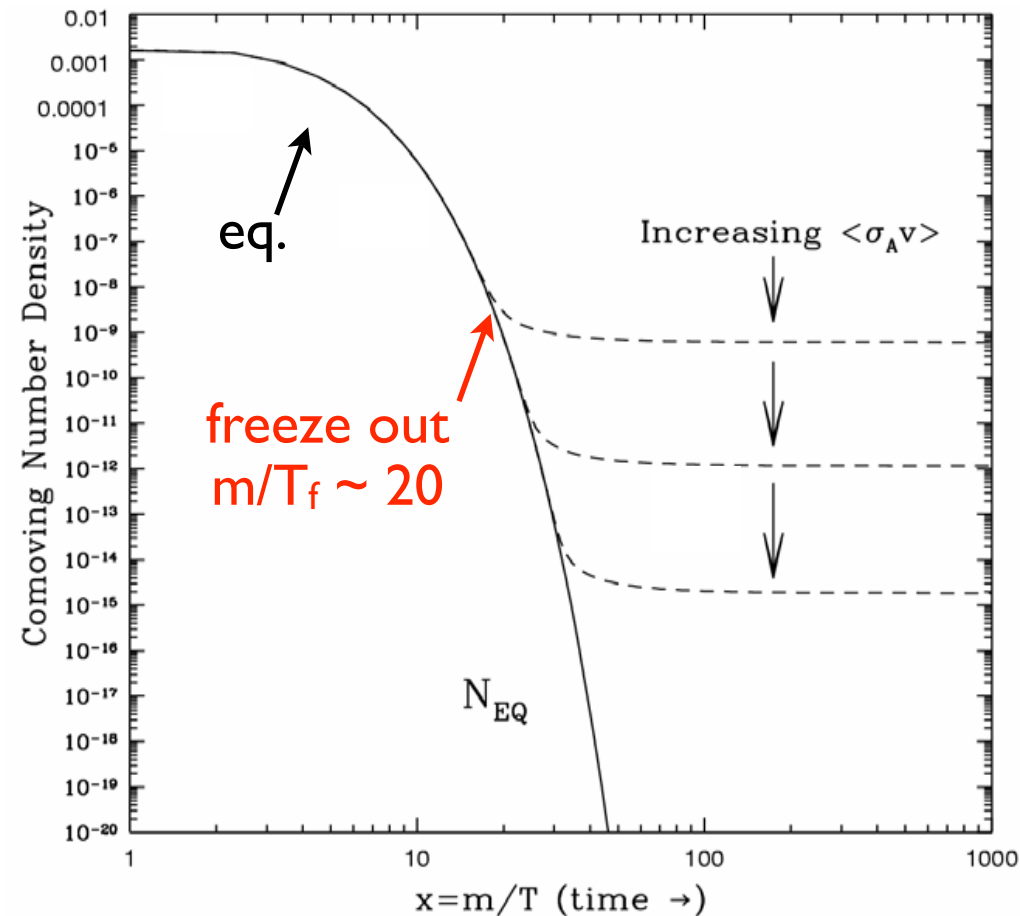
← cold

freeze out

$$\Omega_{\tilde{\chi}_1^0} h^2 = m_{\tilde{\chi}_1^0} Y_{\tilde{\chi}_1^0}^{\text{dec}} s(T_0) h^2 / \rho_c$$

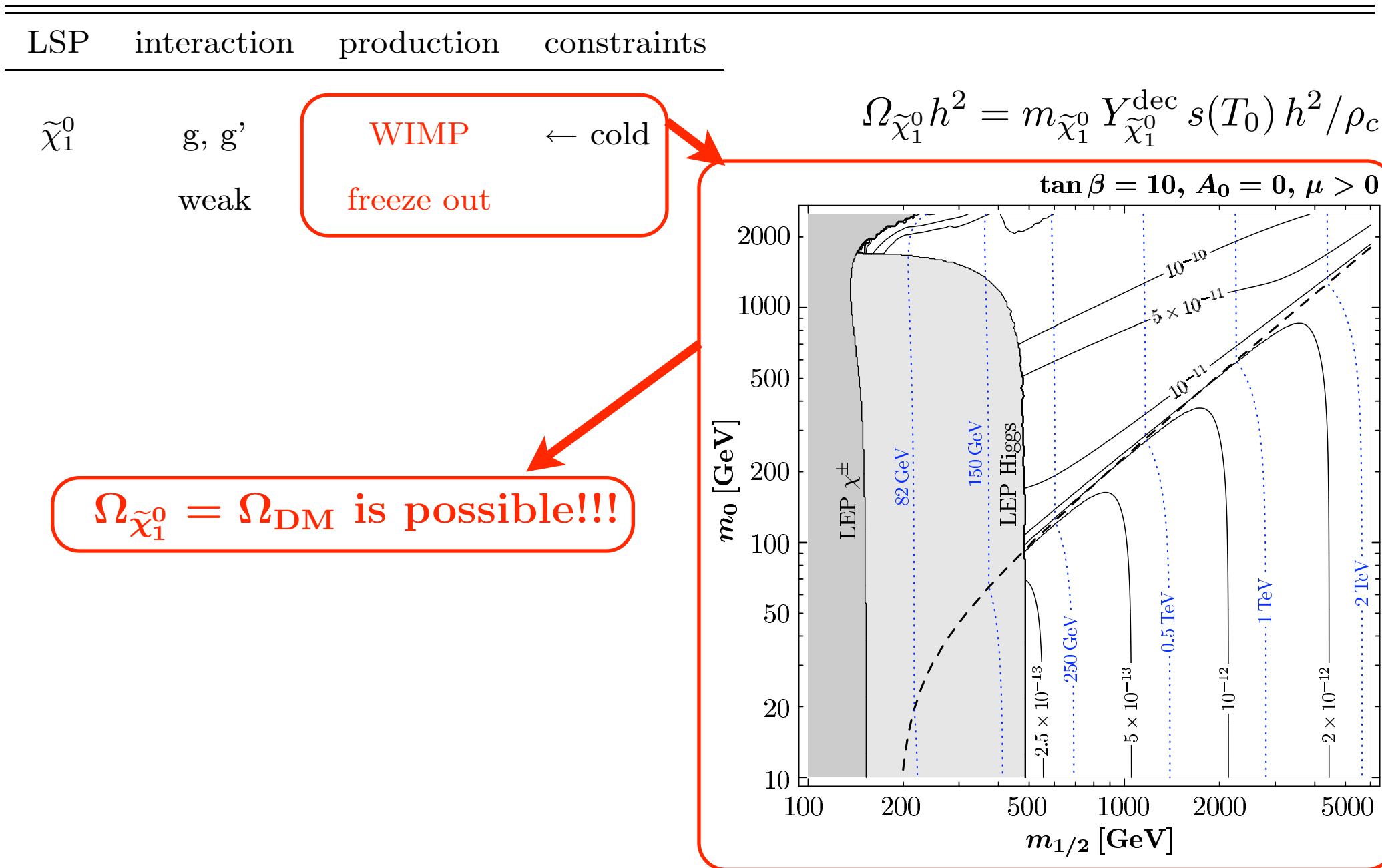
$\Omega_{\tilde{\chi}_1^0} = \Omega_{\text{DM}}$ is possible!!!

Cold Thermal Relic

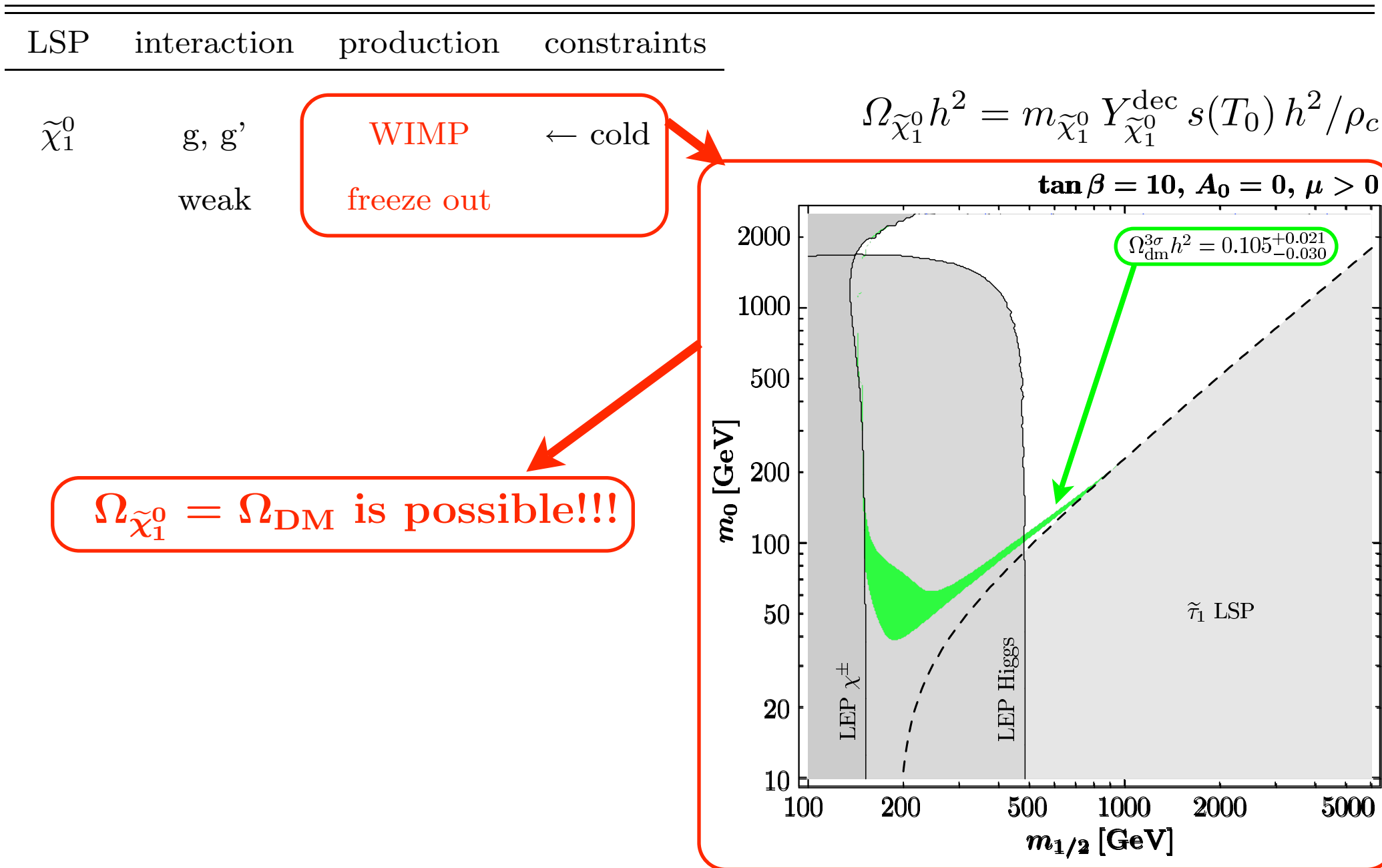


[... , Jungman, Kamionkowski, Griest, '96, ...]

$\tilde{\chi}_1^0$ LSP Dark Matter: Production, Constraints

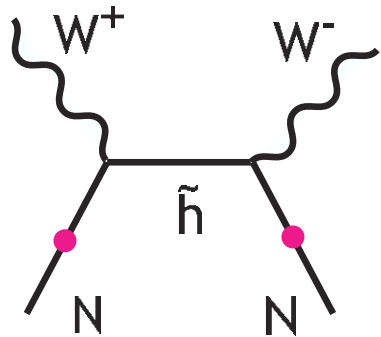


$\tilde{\chi}_1^0$ LSP Dark Matter: Production, Constraints

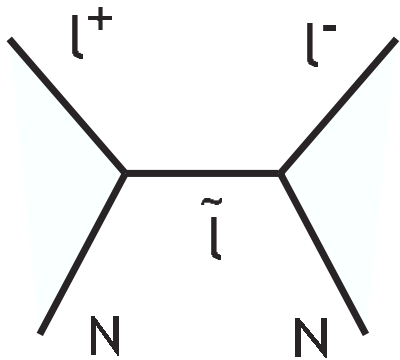


Neutralino LSP Case

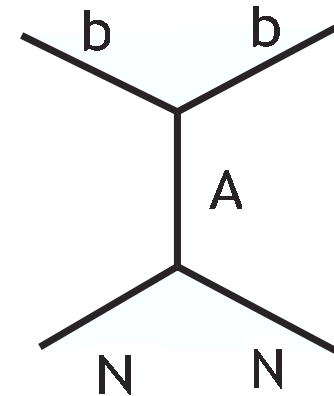
focus point region



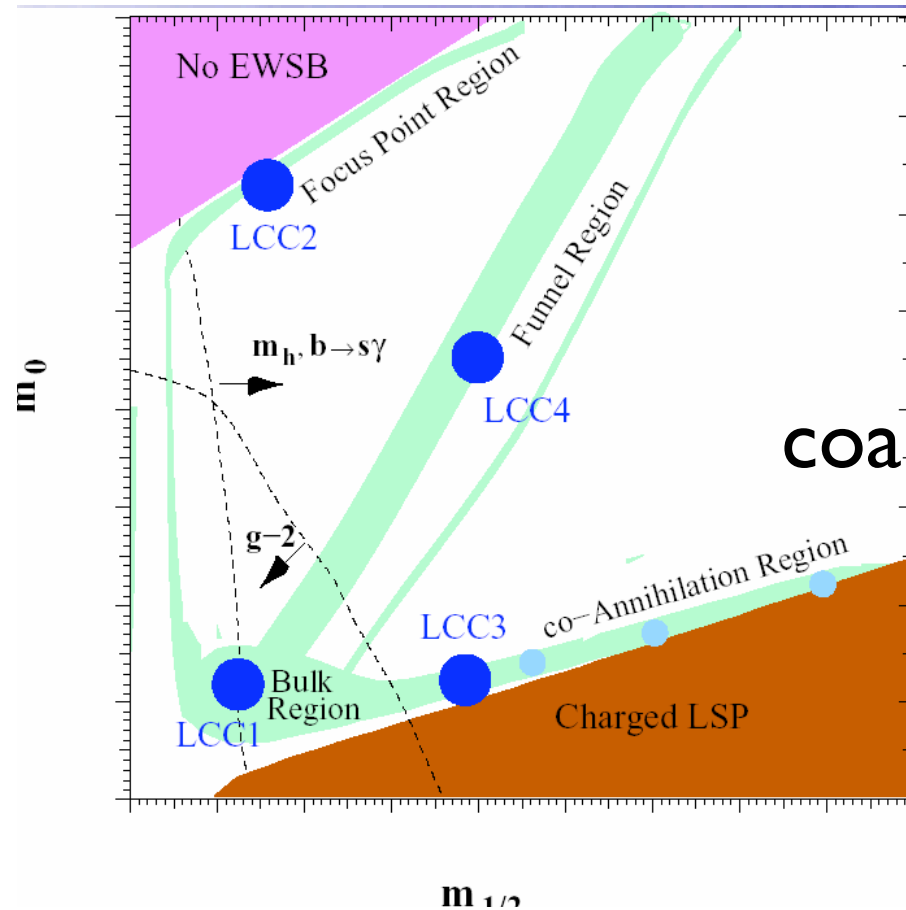
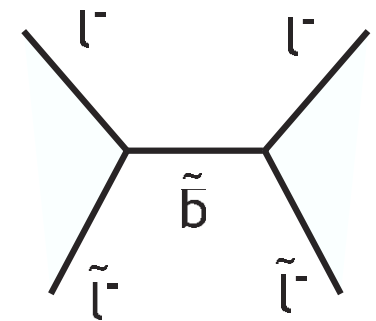
bulk region



annihilation funnel



coannihilation region



[see Baltz, Battaglia, Peskin, Wizansky, '06]

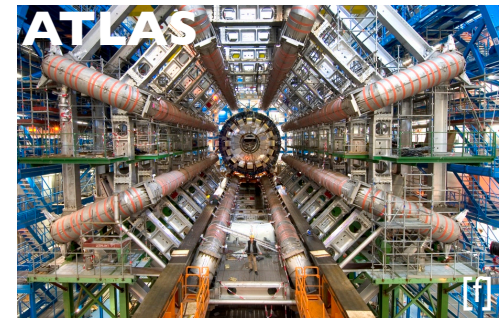
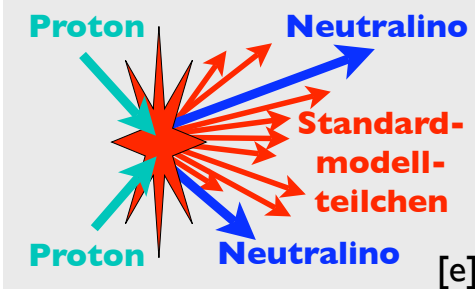
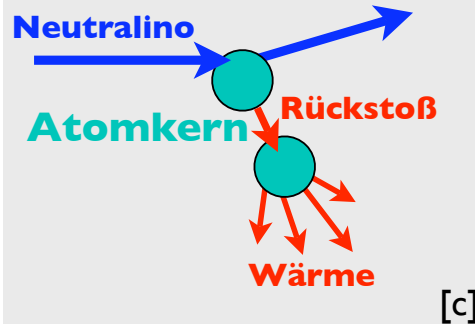
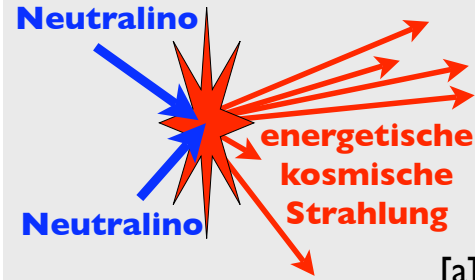
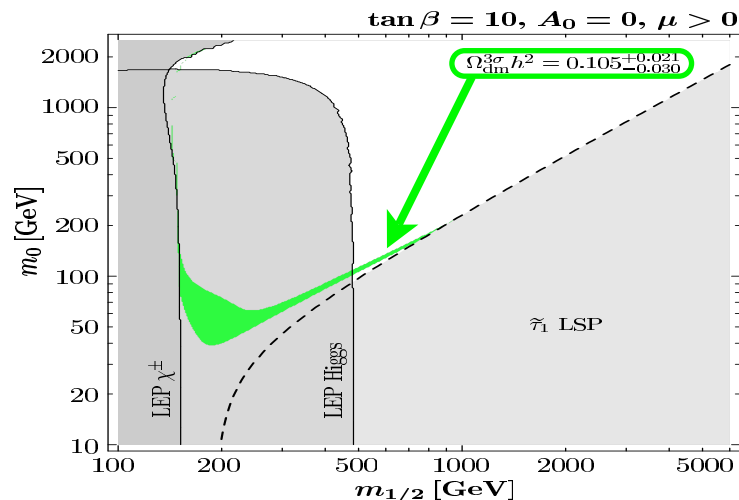
$\tilde{\chi}_1^0$ LSP Dark Matter: Production, Constraints, Experiments

LSP interaction production constraints experiments

$\tilde{\chi}_1^0$ g, g'
weak

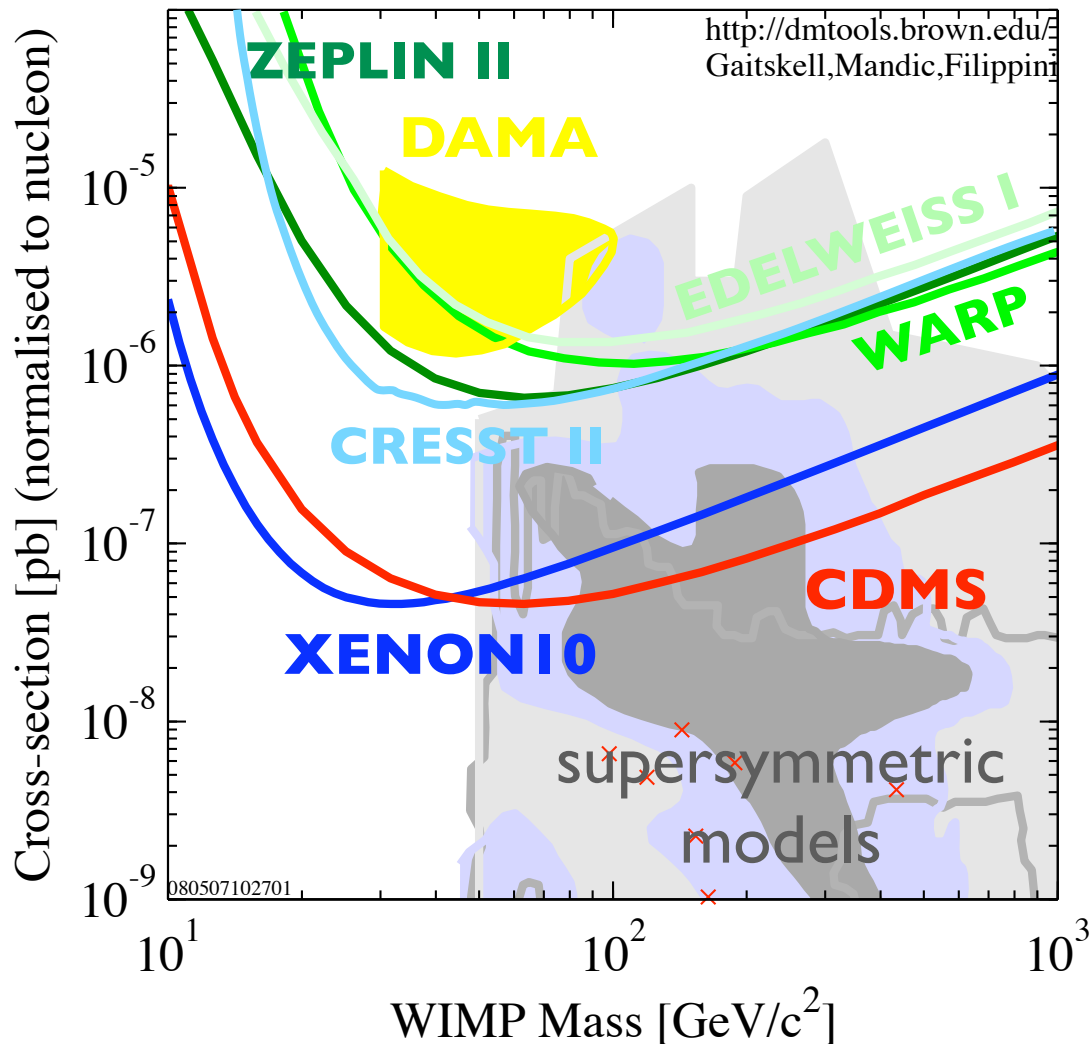
WIMP ← cold
freeze out

$\Omega_{\tilde{\chi}_1^0} = \Omega_{\text{DM}}$ is possible!!!



promising experimental prospects

Upper limits on the neutralino-nucleon interaction strength



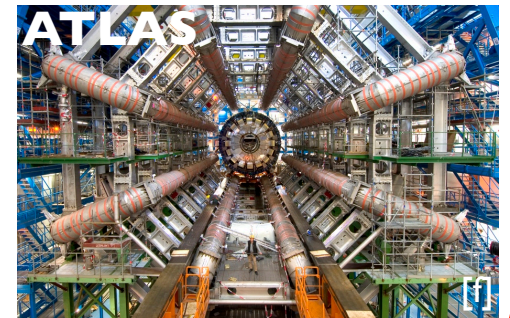
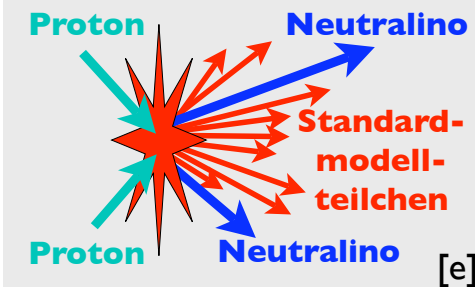
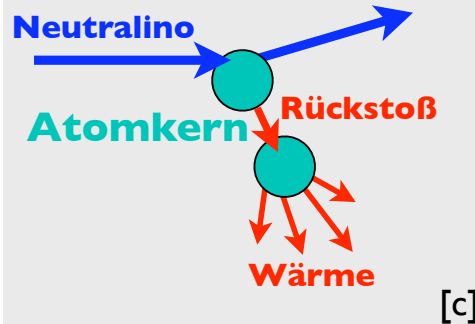
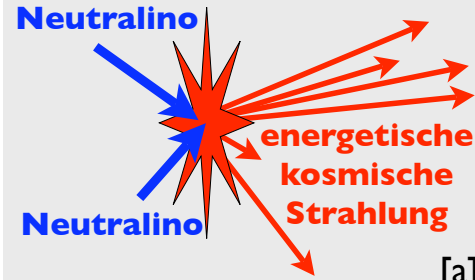
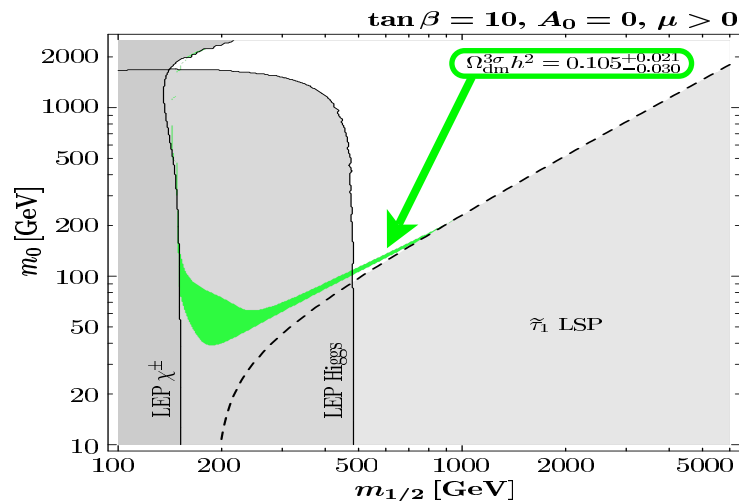
$\tilde{\chi}_1^0$ LSP Dark Matter: Production, Constraints, Experiments

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weak

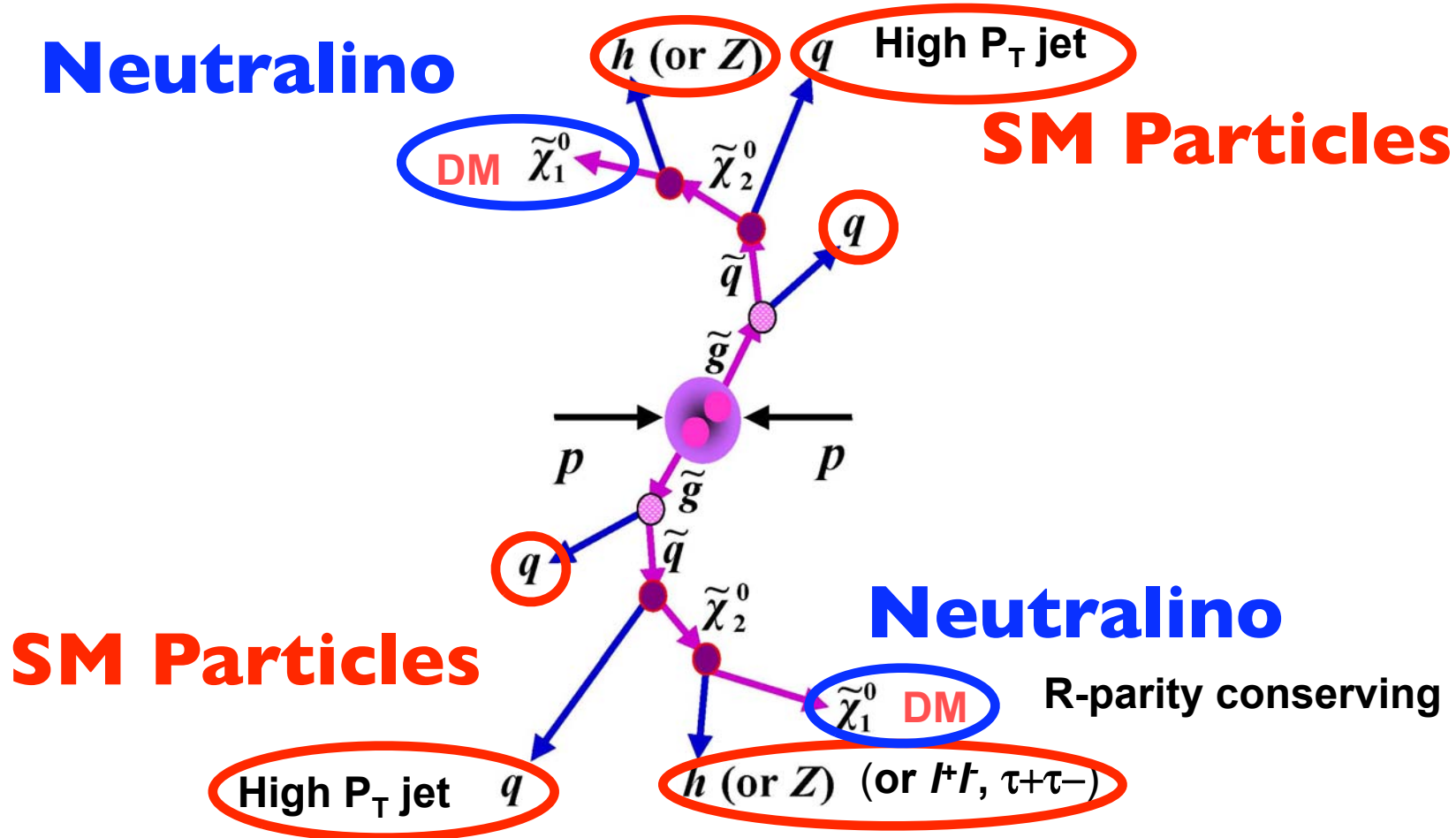
WIMP ← cold
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promising experimental prospects

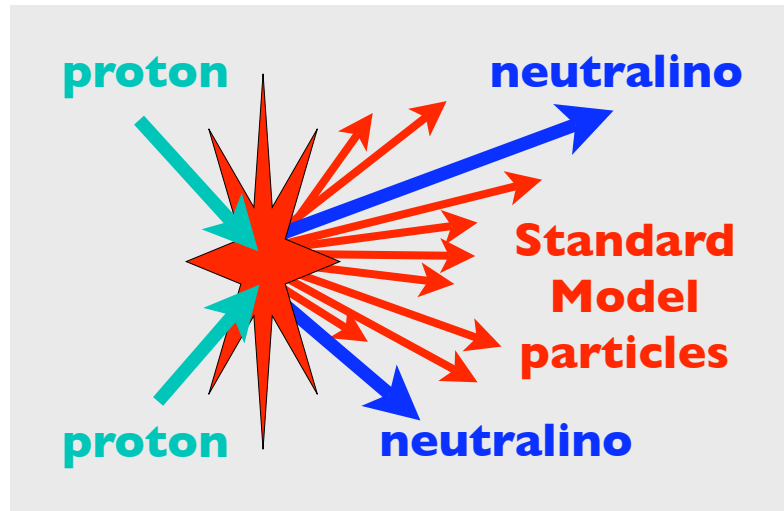
Neutralino DM Production at the LHC



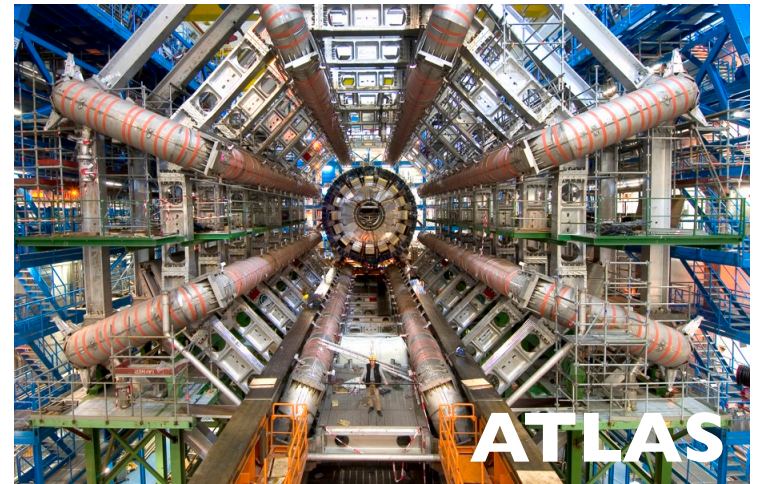
The signal : **jets + leptons + missing E_T**

[from B. Dutta's Talk, SUSY 2007]

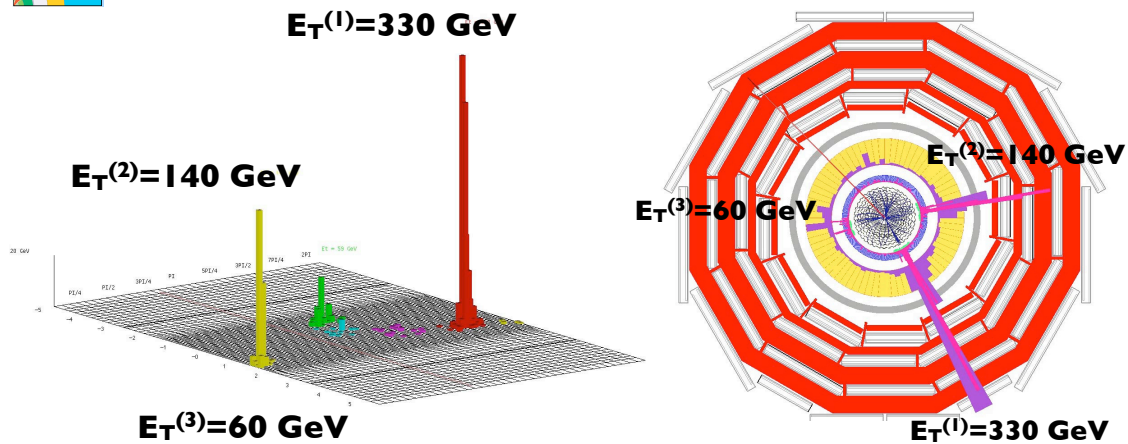
Collider Searches



ongoing
searches
at
Tevatron
pp @ 2 TeV
CDF D0



$E_T^{\text{missing}} = 360 \text{ GeV}$



The signal:
jets + leptons + large E_T^{miss}



Warning



**Things might turn out
to be very different ...**

Other well-motivated candidates



Extremely Weakly Interacting Particles (EWIPs)

Extensions of the Standard Model

Peccei-Quinn Symmetry & Supersymmetry

	Axions $f_a > 10^9 \text{ GeV}$	Axinos $f_a > 10^9 \text{ GeV}$	Gravitinos $M_{\text{Pl}} = 2.4 \times 10^{18} \text{ GeV}$
spin	0	1/2	3/2
mass	$< 10 \text{ meV}$?	eV-TeV
int.	$\propto (p/f_a)^n$	$\propto (p/f_a)^n$	$\propto (p/M_{\text{Pl}})^n$

[Talk by Jihn E. Kim]

Other well-motivated candidates



Extremely Weakly Interacting Particles (EWIPs)

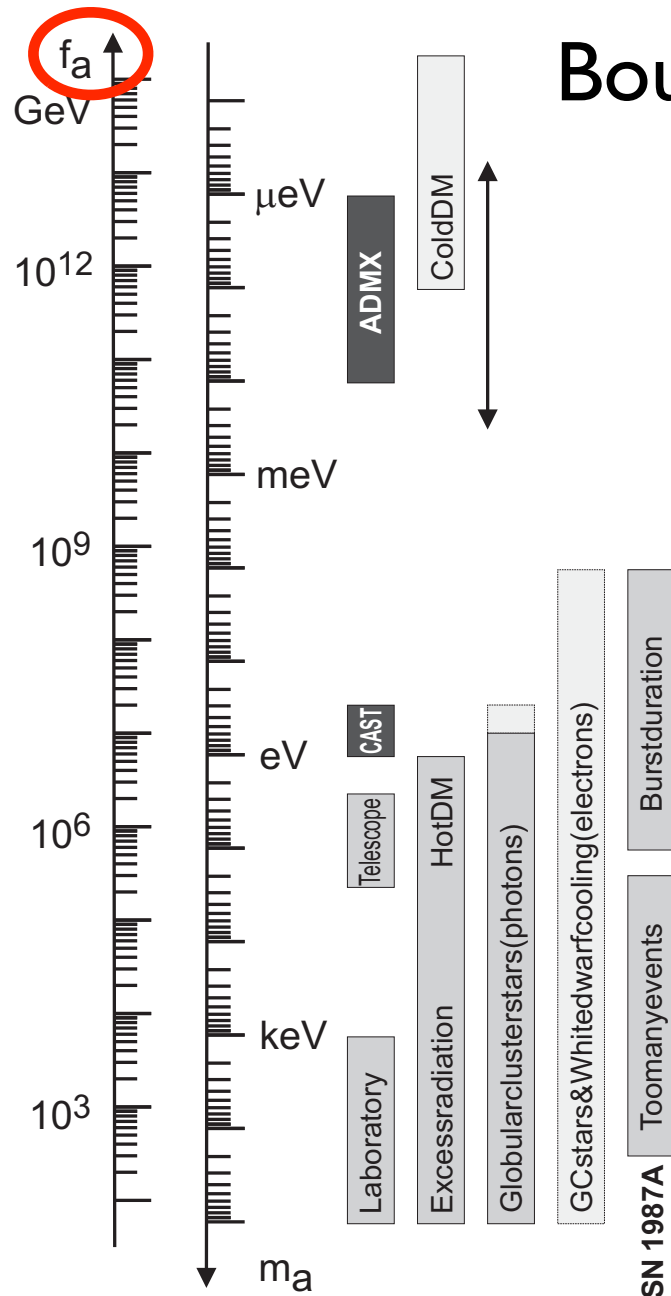
Extensions of the Standard Model

Peccei-Quinn Symmetry & Supersymmetry

	Axions $f_a > 10^9 \text{ GeV}$	Axinos $f_a > 10^9 \text{ GeV}$	Gravitinos $M_{\text{Pl}} = 2.4 \times 10^{18} \text{ GeV}$
spin	0	1/2	3/2
mass	$< 10 \text{ meV}$?	eV-TeV
int.	$\propto (p/f_a)^n$	$\propto (p/f_a)^n$	$\propto (p/M_{\text{Pl}})^n$

[Talk by Jihn E. Kim]

Bounds on the Peccei-Quinn Scale

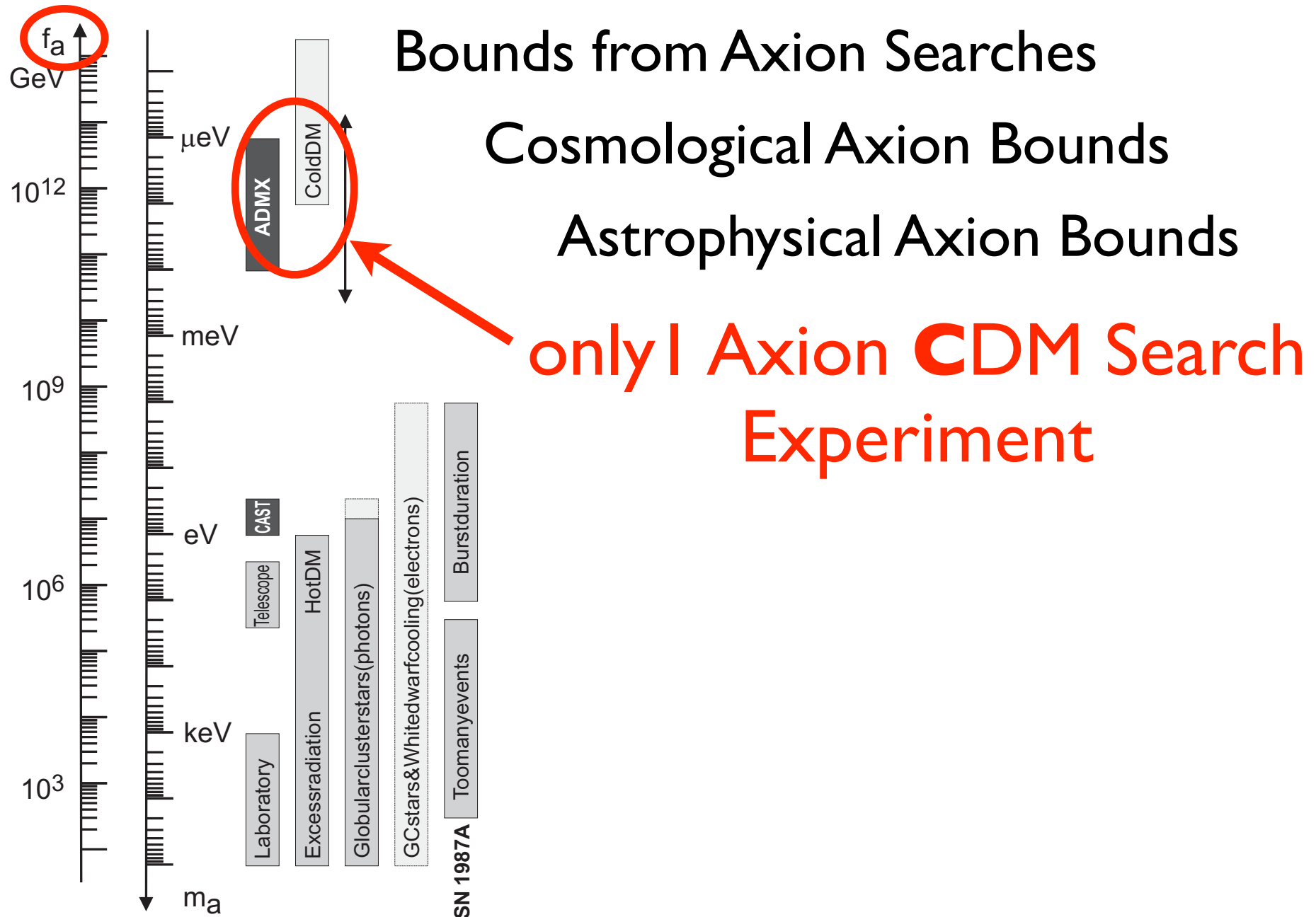


Bounds from Axion Searches

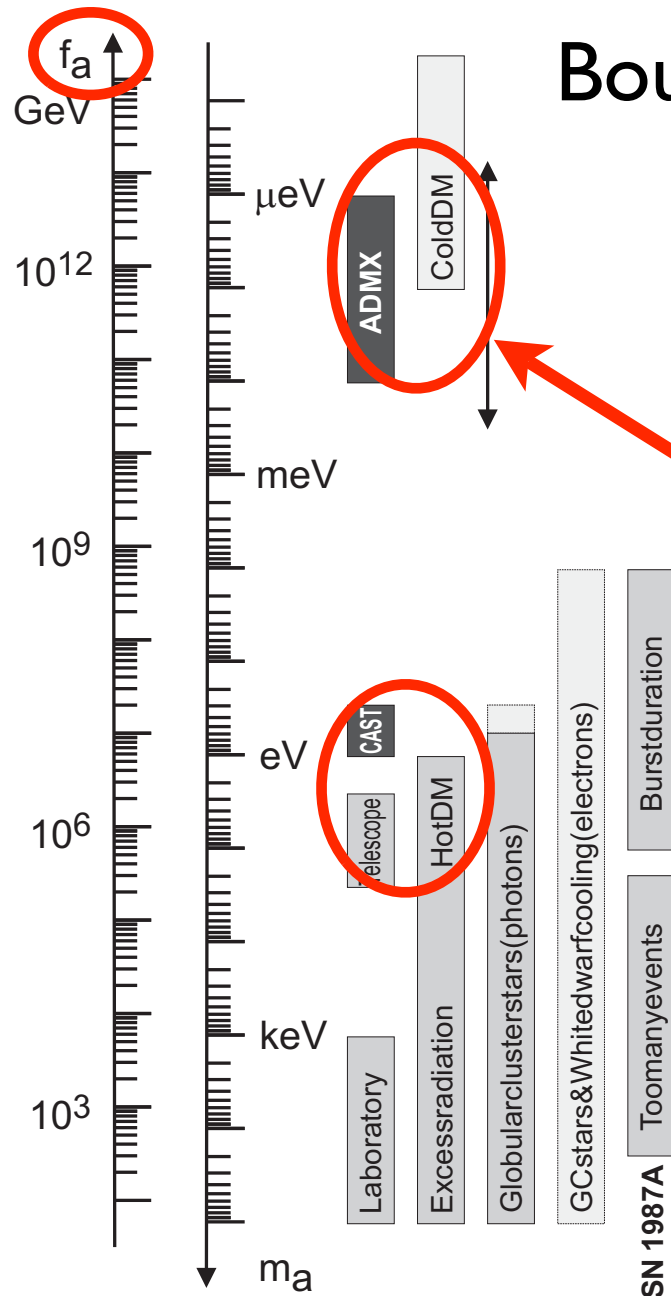
Cosmological Axion Bounds

Astrophysical Axion Bounds

Bounds on the Peccei-Quinn Scale



Bounds on the Peccei-Quinn Scale



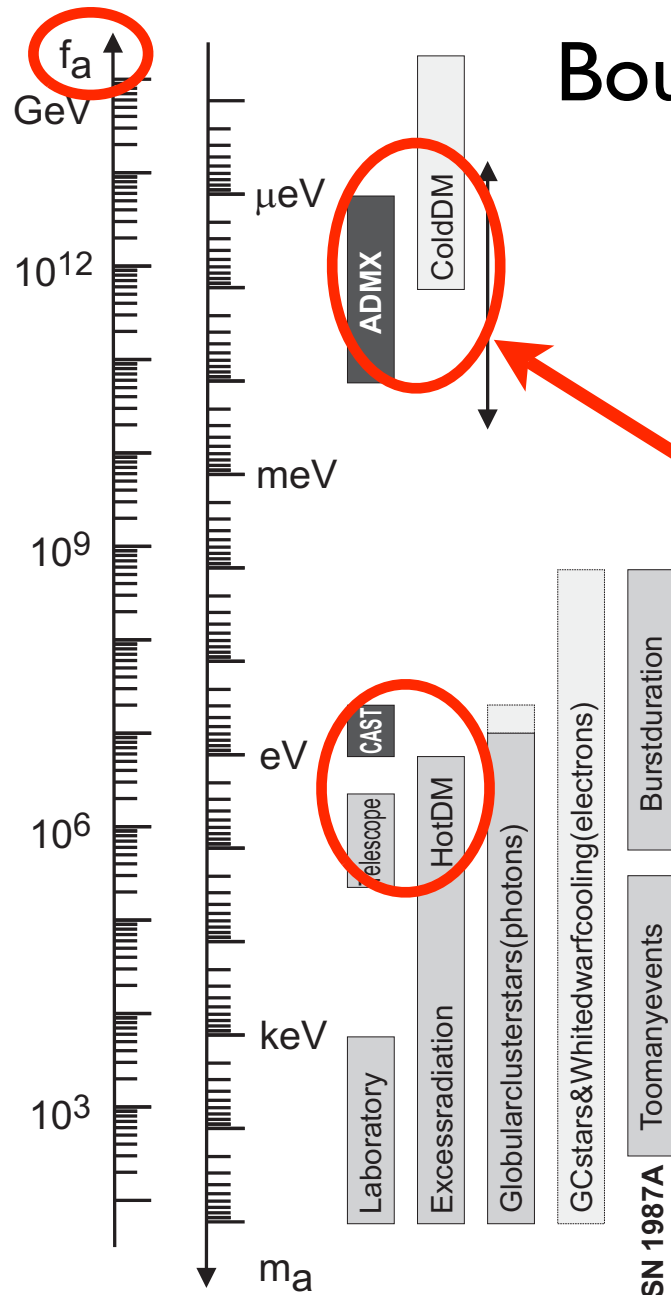
Bounds from Axion Searches

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Astrophysical Axion Bounds

only 1 Axion CDM Search Experiment

Bounds on the Peccei-Quinn Scale



Bounds from Axion Searches

Cosmological Axion Bounds

Astrophysical Axion Bounds

only 1 Axion CDM Search Experiment

Axion DM and
Neutralino DM
might coexist!

Extremely Weakly Interacting Particles (EWIPs)

Extensions of the Standard Model

Peccei-Quinn Symmetry & Supersymmetry

	Axions $f_a > 10^9 \text{ GeV}$	Axinos $f_a > 10^9 \text{ GeV}$	Gravitinos $M_{\text{Pl}} = 2.4 \times 10^{18} \text{ GeV}$
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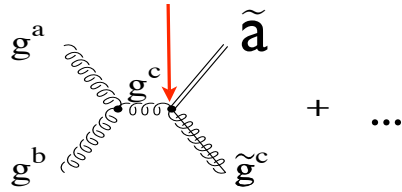
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Thermal Axino Production



inflation

radiation dominated

mat. dom.

Λ dom.

$$\rho_{\text{rad}} \propto a^{-4}$$

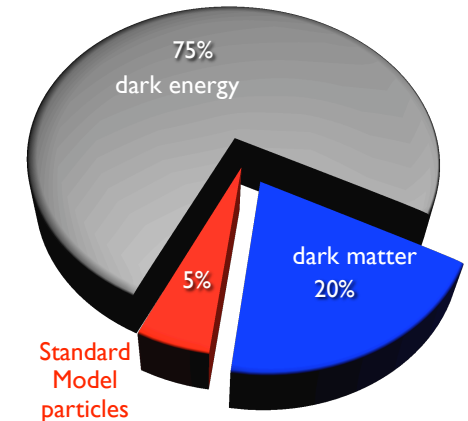
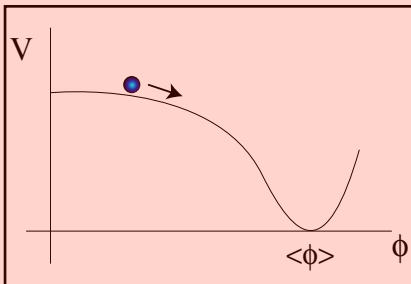
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slow
roll

reheat
phase

$$\rho_{\phi} \propto a^0$$



1 s

100.000 y

$t_0 = 14 \text{ Gy}$

t

$T_R = ?$

10 GeV

1 MeV

1 eV

$T_0 = 2.73 \text{ K}$

T

reheating
temp.

WIMP
freeze
out

BBN
CvB

CMB

LHC

a

Thermal Production of Axino Dark Matter in the Early Universe

Axino Number Density for $f_a > T_D > T_R \gtrsim 10^4$ GeV

- Boltzmann equation: time evolution of axino density $n_{\tilde{a}}$ in the thermal bath

$$\frac{dn_{\tilde{a}}}{dt} + 3Hn_{\tilde{a}} = C_{\tilde{a}} = \int d^3p \frac{d\Gamma_{\tilde{a}}}{d^3p} \quad \leftarrow \quad \text{generation of } \tilde{a} - \text{annihilation of } \tilde{a}$$

- collision term for $a(p_1) + b(p_2) \rightarrow c(p_3) + \tilde{a}(p)$: $(C_{a+b \rightarrow c+\tilde{a}} \in C_{\tilde{a}})$

$$C_{a+b \rightarrow c+\tilde{a}} = \int \frac{d^3p}{(2\pi)^3 2E} \int \left[\prod_{i=1}^3 \frac{d^3p_i}{(2\pi)^3 2E_i} \right] (2\pi)^4 \delta^4(p_1 + p_2 - p_3 - p) \\ \times \left[|M_{a+b \rightarrow c+\tilde{a}}|^2 f_a f_b (1 \pm f_c) (1 - f_{\tilde{a}}) - |M_{c+\tilde{a} \rightarrow a+b}|^2 f_c f_{\tilde{a}} (1 \pm f_a) (1 \pm f_b) \right]$$

- phase space densities: $f_i \rightarrow$ number densities: $n_i = \int \frac{d^3p_i}{(2\pi)^3 2E_i} g_i f_i(E_i)$

$$a, b, \text{ and } c: f_i = f_i^{\text{eq}} = f_{B/F} = \frac{1}{\exp(E_i/T) \mp 1} \quad , \quad \text{axino: } f_{\tilde{a}} \approx 0$$

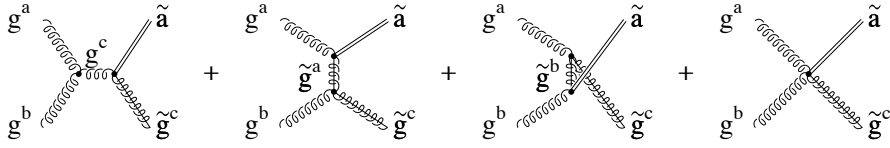
Axino Interactions \longleftarrow Hadronic (KSVZ) Axion Models

- axino–gluino–gluon interaction:

$$\mathcal{L}_{\tilde{a}\tilde{g}g} = i \frac{\alpha_s}{16\pi(f_a/N)} \bar{\tilde{a}} \gamma_5 [\gamma^\mu, \gamma^\nu] \tilde{g}^a G_{\mu\nu}^a$$

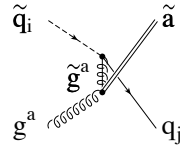
Thermal Axino Production in SUSY QCD

- A: $g^a + g^b \rightarrow \tilde{g}^c + \tilde{a}$



- B: $g^a + \tilde{g}^b \rightarrow g^c + \tilde{a}$ (crossing of A)

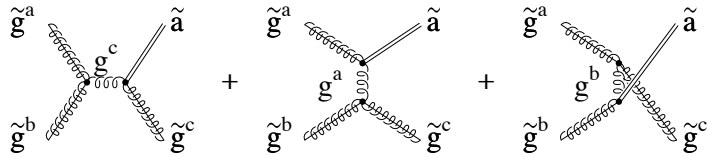
- C: $\tilde{q}_i + g^a \rightarrow \tilde{q}_j + \tilde{a}$



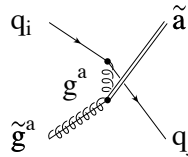
- D: $g^a + q_i \rightarrow \tilde{q}_j + \tilde{a}$ (crossing of C)

- E: $\bar{\tilde{q}}_i + q_j \rightarrow g^a + \tilde{a}$ (crossing of C)

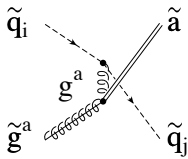
- F: $\tilde{g}^a + \tilde{g}^b \rightarrow \tilde{g}^c + \tilde{a}$



- G: $q_i + \tilde{g}^a \rightarrow q_j + \tilde{a}$



- H: $\tilde{q}_i + \tilde{g}^a \rightarrow \tilde{q}_j + \tilde{a}$



- I: $q_i + \bar{q}_j \rightarrow \tilde{g}^a + \tilde{a}$ (crossing of G)

- J: $\tilde{q}_i + \bar{\tilde{q}}_j \rightarrow \tilde{g}^a + \tilde{a}$ (crossing of H)

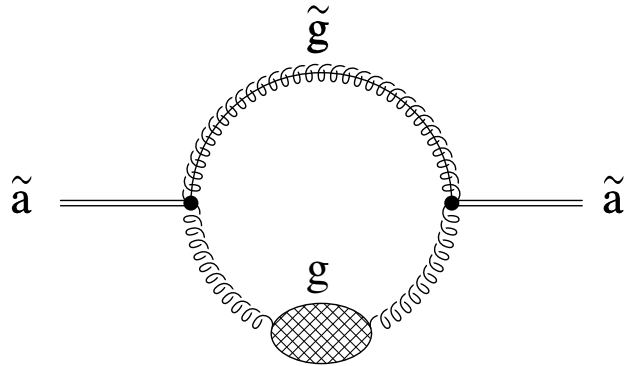
Axino Production in $2 \rightarrow 2$ Processes

	process i	$ \mathcal{M}_i ^2 / \frac{g^6}{128\pi^4(f_a/N)^2}$
A	$g^a + g^b \rightarrow \tilde{g}^c + \tilde{a}$	$4(s + 2t + 2\frac{t^2}{s}) f^{abc} ^2$
B	$g^a + \tilde{g}^b \rightarrow g^c + \tilde{a}$	$-4(t + 2s + 2\frac{s^2}{t}) f^{abc} ^2$
C	$\tilde{q}_i + g^a \rightarrow q_j + \tilde{a}$	$2s T_{ji}^a ^2$
D	$g^a + q_i \rightarrow \tilde{q}_j + \tilde{a}$	$-2t T_{ji}^a ^2$
E	$\bar{\tilde{q}}_i + q_j \rightarrow g^a + \tilde{a}$	$-2t T_{ji}^a ^2$
F	$\tilde{g}^a + \tilde{g}^b \rightarrow \tilde{g}^c + \tilde{a}$	$-8\frac{(s^2 + st + t^2)^2}{st(s+t)} f^{abc} ^2$
G	$q_i + \tilde{g}^a \rightarrow q_j + \tilde{a}$	$-4(s + \frac{s^2}{t}) T_{ji}^a ^2$
H	$\tilde{q}_i + \tilde{g}^a \rightarrow \tilde{q}_j + \tilde{a}$	$-2(\frac{t}{2} + 2s + 2\frac{s^2}{t}) T_{ji}^a ^2$
I	$q_i + \bar{q}_j \rightarrow \tilde{g}^a + \tilde{a}$	$-4(t + \frac{t^2}{s}) T_{ji}^a ^2$
J	$\tilde{q}_i + \bar{\tilde{q}}_j \rightarrow \tilde{g}^a + \tilde{a}$	$2(\frac{s}{2} + 2t + 2\frac{t^2}{s}) T_{ji}^a ^2$

B, F, G, & H: Logarithmic IR Singularity

- Separation of Scales: $gT \ll \Lambda \ll T \leftarrow g \ll 1$ [Braaten, Yuan, 1991]

- Soft Part: Axino Self-Energy

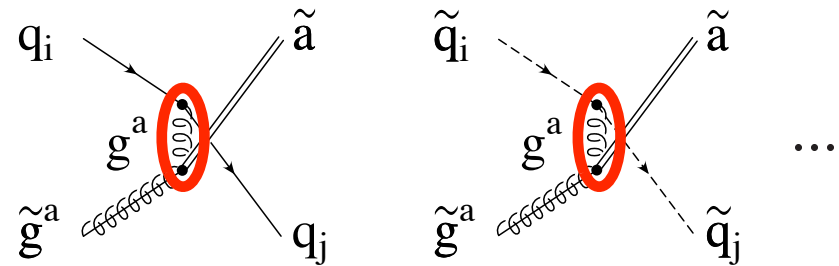


→ eff. HTL resummed propagator

$$E \left. \frac{d\Gamma_{\tilde{a}}}{d^3p} \right|_{\text{soft}} = -f_F(E) \frac{\text{Im}\Sigma(E + i\epsilon, \mathbf{p})}{(2\pi)^3} \Big|_{|\mathbf{p}_1 - \mathbf{p}_3| < \Lambda}$$

$$= A_{\text{soft}} + B \ln \left[\frac{\Lambda}{gT} \right]$$

- Hard Part: Relativ. Kin. Theory



→ bare propagator

$$E \left. \frac{d\Gamma_{\tilde{a}}}{d^3p} \right|_{\text{hard}} = \frac{1}{2} \int \prod_{i=1}^3 \dots |M|^2 \Theta(|\mathbf{p}_1 - \mathbf{p}_3| - \Lambda)$$

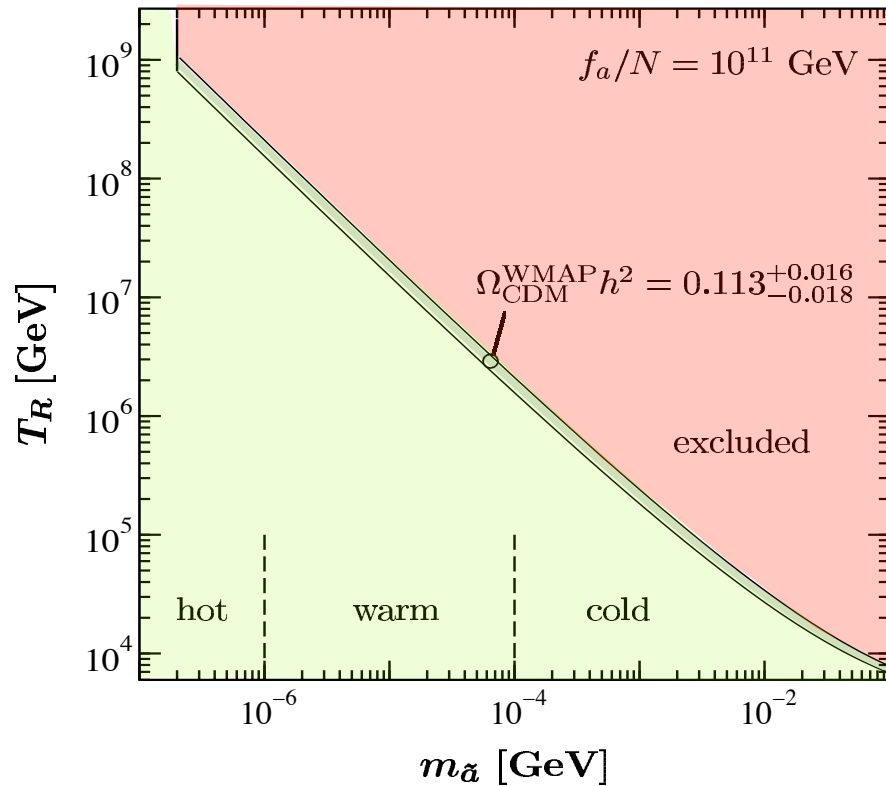
$$= A_{\text{hard}} + B \ln \left[\frac{T}{\Lambda} \right]$$

- Thermal Production Rate: * complete to LO in g , * finite , * indep. of Λ

$$E \left. \frac{d\Gamma_{\tilde{G}}}{d^3p} \right|_{\text{LO in } g} = E \left. \frac{d\Gamma_{\tilde{G}}}{d^3p} \right|_{\text{soft}} + E \left. \frac{d\Gamma_{\tilde{G}}}{d^3p} \right|_{\text{hard}} = A_{\text{soft}} + A_{\text{hard}} + B \ln \left[\frac{1}{g} \right]$$

Axino LSP Case

Thermal \tilde{a} Production



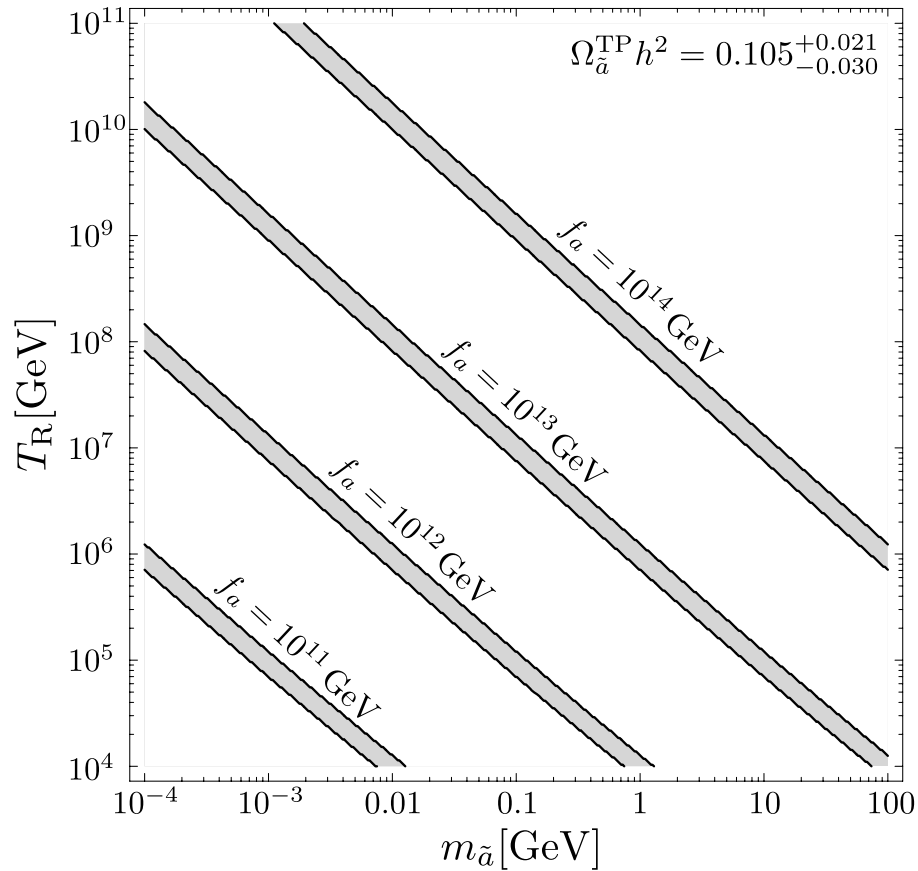
[Brandenburg, FDS, '04]

see also [Covi et al., '01]

and [Strumia, '10]

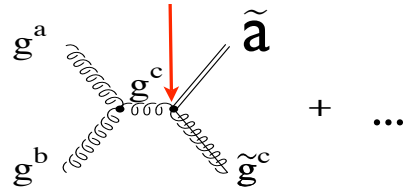
Axino LSP Case

Thermal \tilde{a} Production



[Freitas, FDS, Tajuddin, Wyler, '09]

Thermal Axino Production



inflation

radiation dominated

mat. dom.

Λ dom.

$$\rho_{\text{rad}} \propto a^{-4}$$

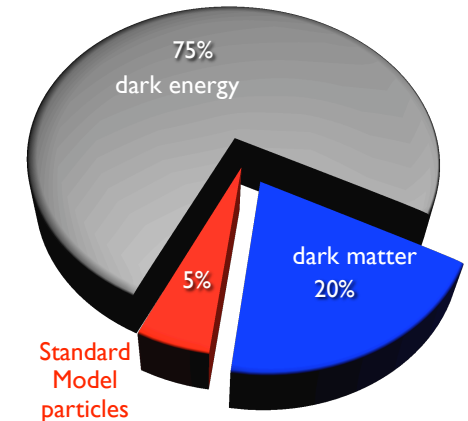
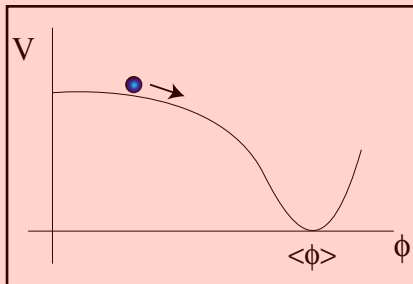
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slow
roll

reheat
phase

$$\rho_{\phi} \propto a^0$$



1 s

100.000 y

$t_0 = 14 \text{ Gy}$

t

$T_R = ?$

10 GeV

1 MeV

1 eV

$T_0 = 2.73 \text{ K}$

T

reheating
temp.

WIMP
freeze
out

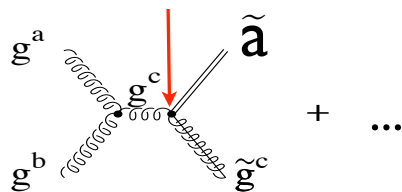
BBN
CvB

CMB

LHC

a

Thermal Axino Production



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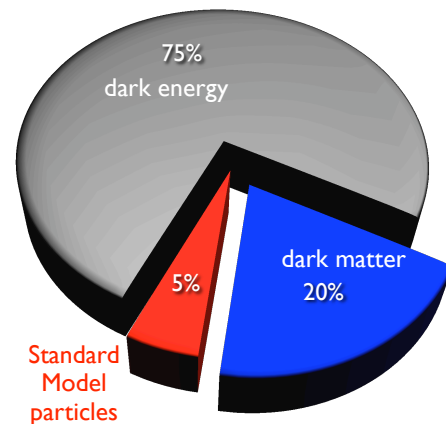
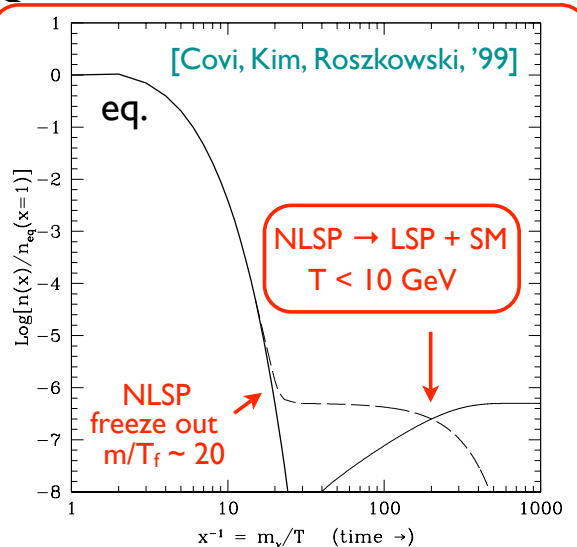
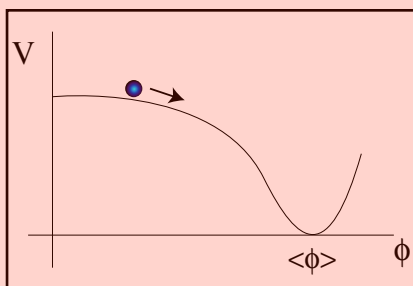
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$T_R = ?$
reheating temp.

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WIMP freeze out

1 MeV
BBN
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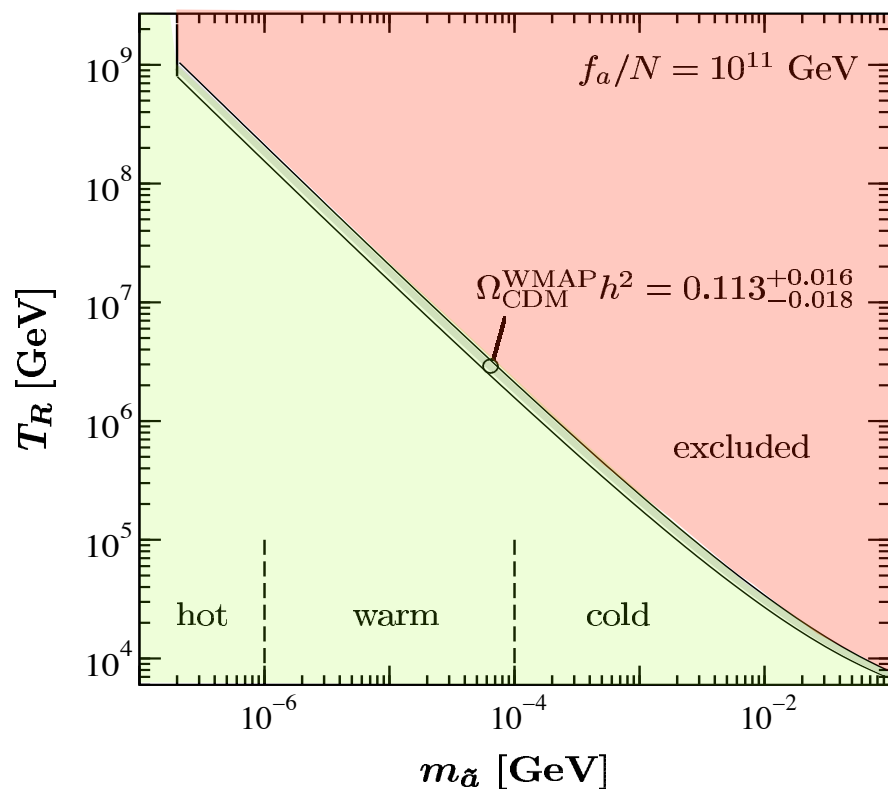
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LHC

t
 T
 a

Axino LSP Case

Thermal \tilde{a} Production

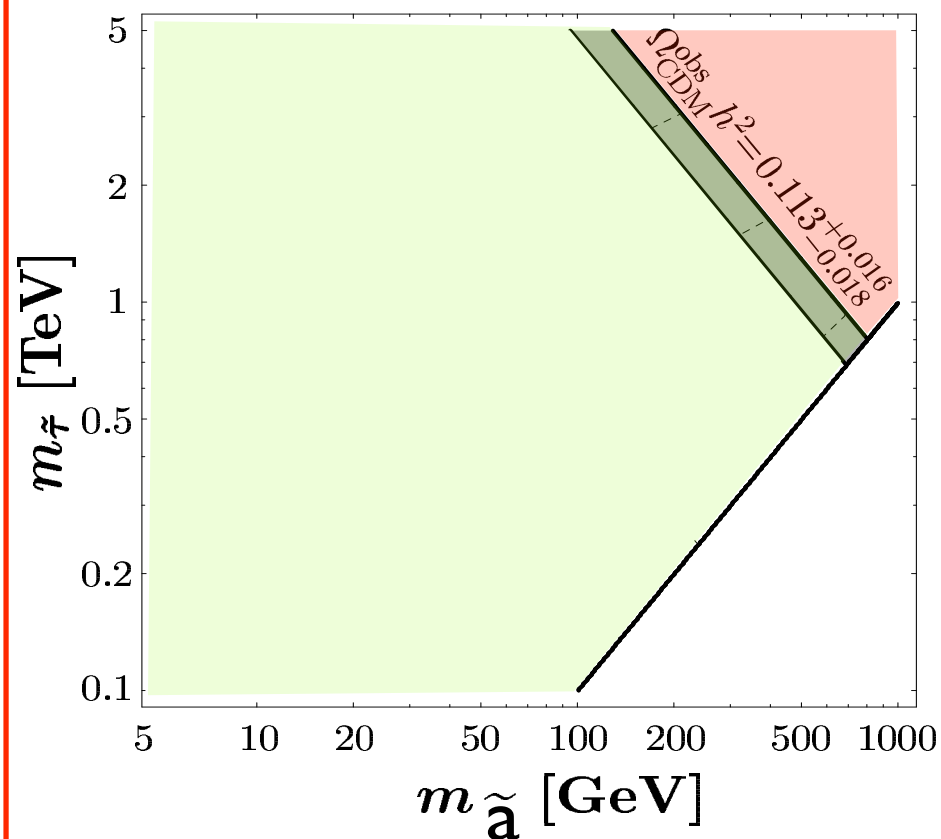


[Brandenburg, FDS, '04]

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$\tilde{\tau}$ NLSP $\rightarrow \tilde{a} + \tau$



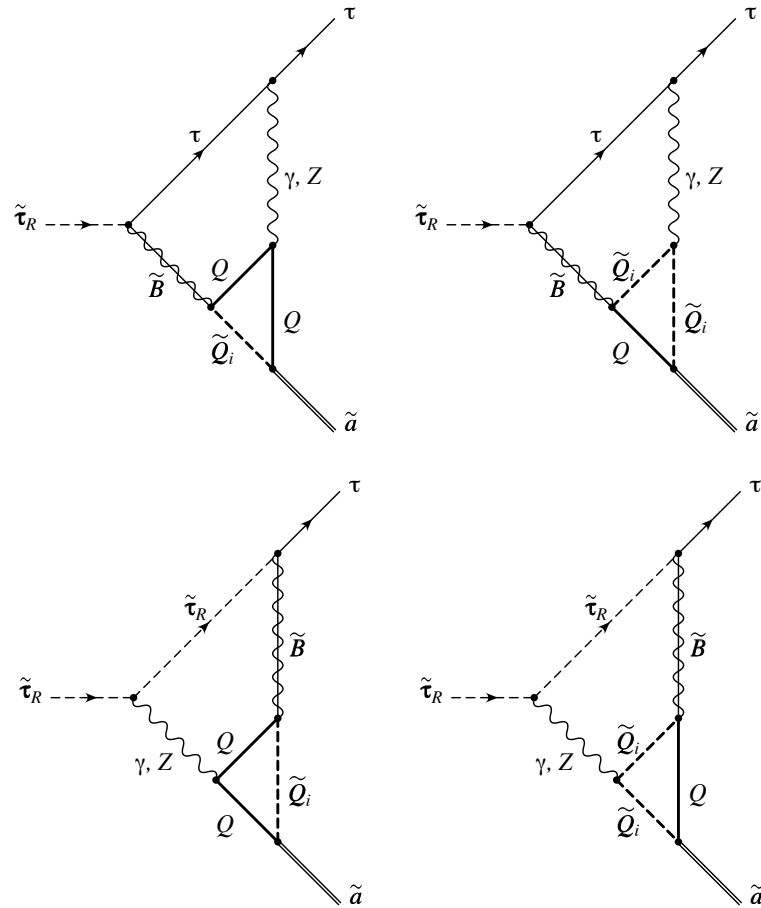
**Probing axinos
experimentally ???**



If we are lucky ...

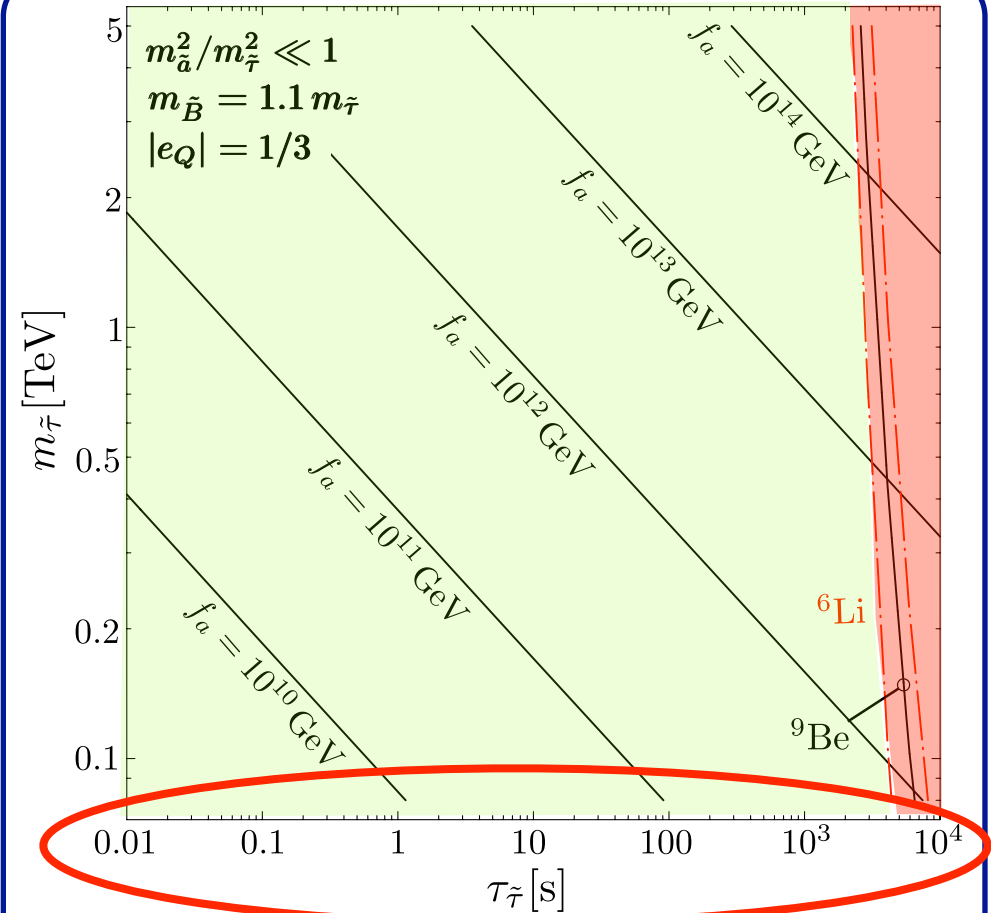
Stau Decays into Axinos

BBN



$$W_{PQ} = y\Phi Q_1 Q_2 \quad M_{\tilde{Q}_{1,2}} = M_Q = y\langle\phi\rangle = yf_a/\sqrt{2}$$

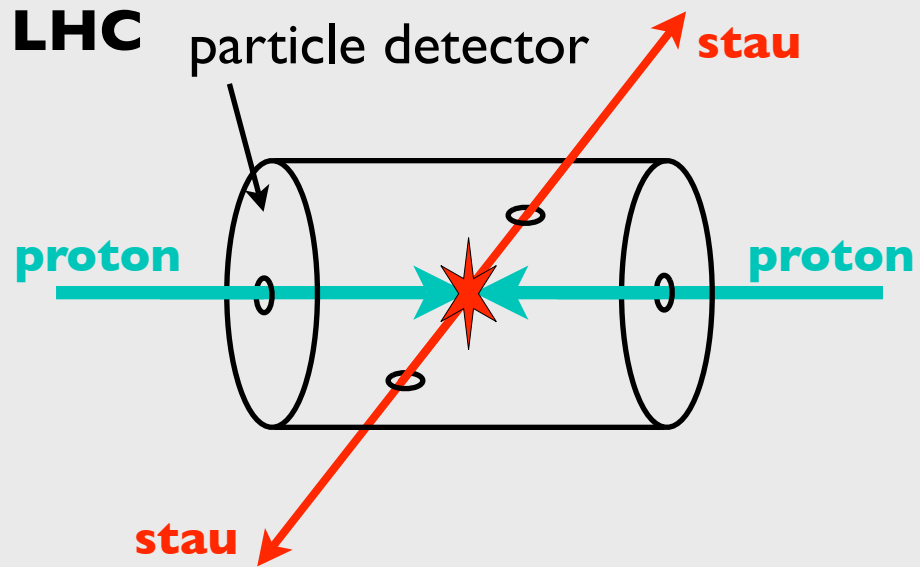
Chiral multiplet	U(1) _{PQ}	(SU(3) _c , SU(2) _L) _Y
$\Phi = \phi + \sqrt{2}\chi\theta + F_\phi\theta\theta$	+1	(1, 1) ₀
$Q_1 = \tilde{Q}_1 + \sqrt{2}q_1\theta + F_1\theta\theta$	-1/2	(3, 1) _{+e_Q}
$Q_2 = \tilde{Q}_2 + \sqrt{2}q_2\theta + F_2\theta\theta$	-1/2	(3*, 1) _{-e_Q}



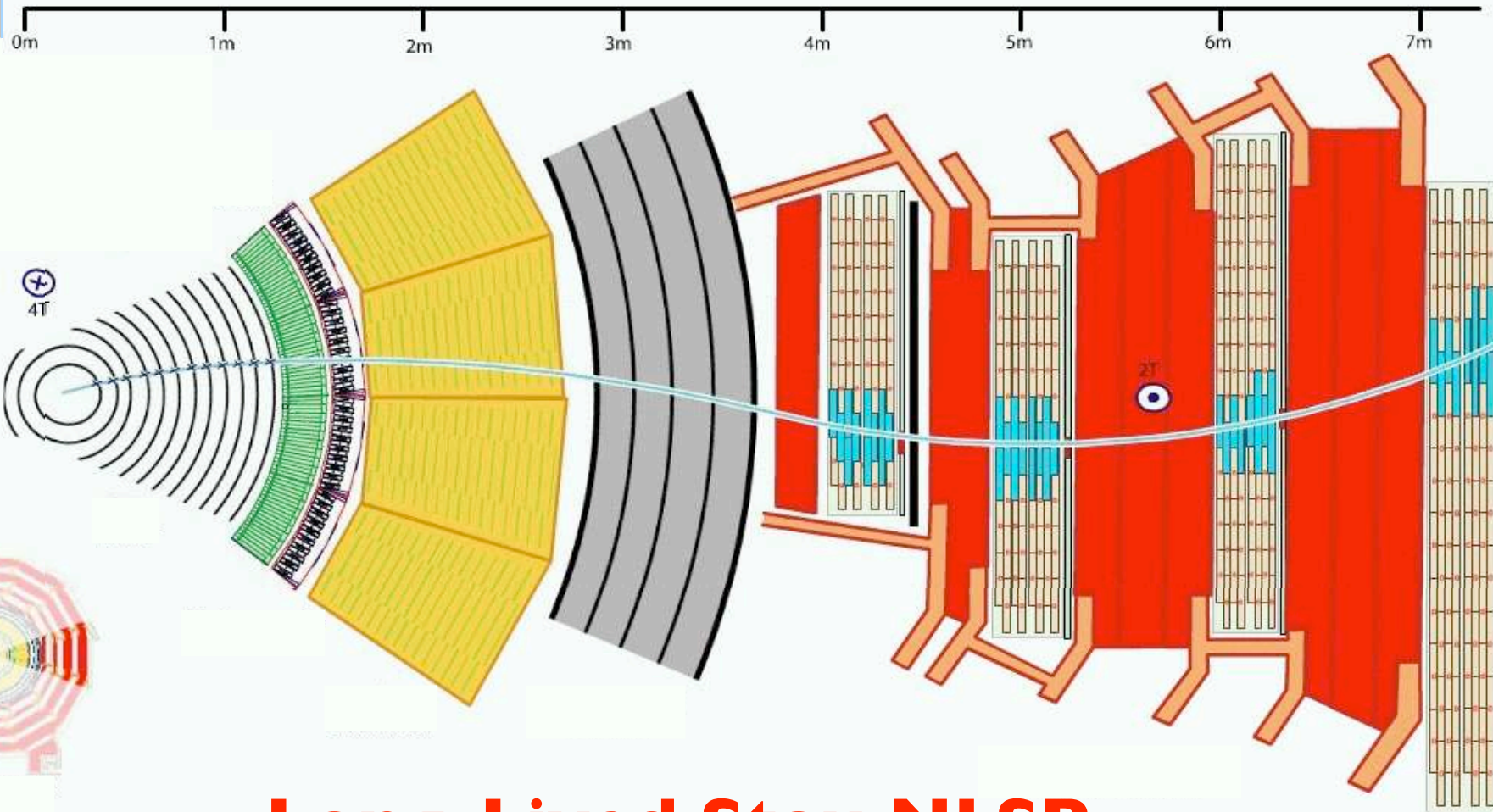
$$\Gamma_{\text{tot}}^{\tilde{\tau}_R} \approx \Gamma(\tilde{\tau}_R \rightarrow \tau \tilde{a})_{\text{LL}}$$

$$= \frac{81\alpha^4 e_Q^4}{128\pi^5 \cos^8 \theta_W} \frac{m_{\tilde{\tau}} m_{\tilde{B}}^2}{f_a^2} \left(1 - \frac{m_{\tilde{a}}^2}{m_{\tilde{\tau}}^2}\right)^2 \ln^2\left(\frac{yf_a}{\sqrt{2}m_{\tilde{\tau}}}\right)$$

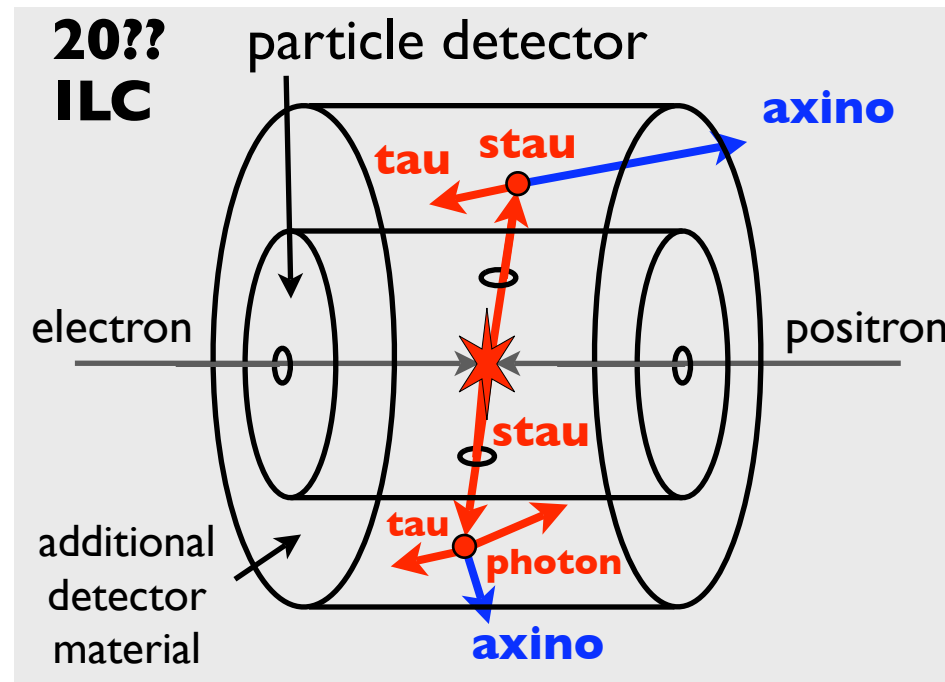
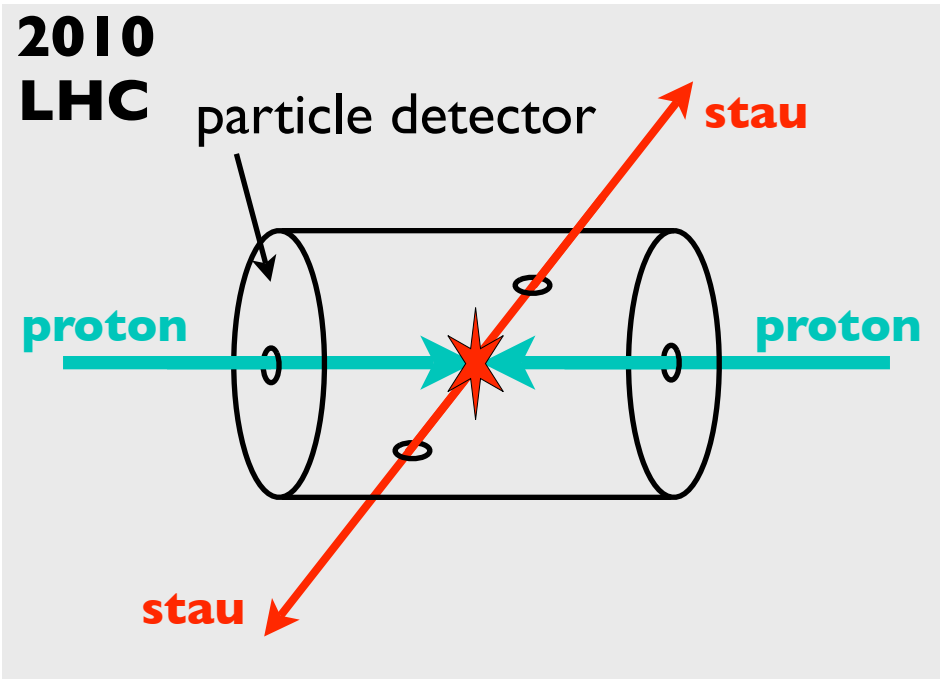
**2010
LHC**



“Stable” Charged Massive Particle @ LHC



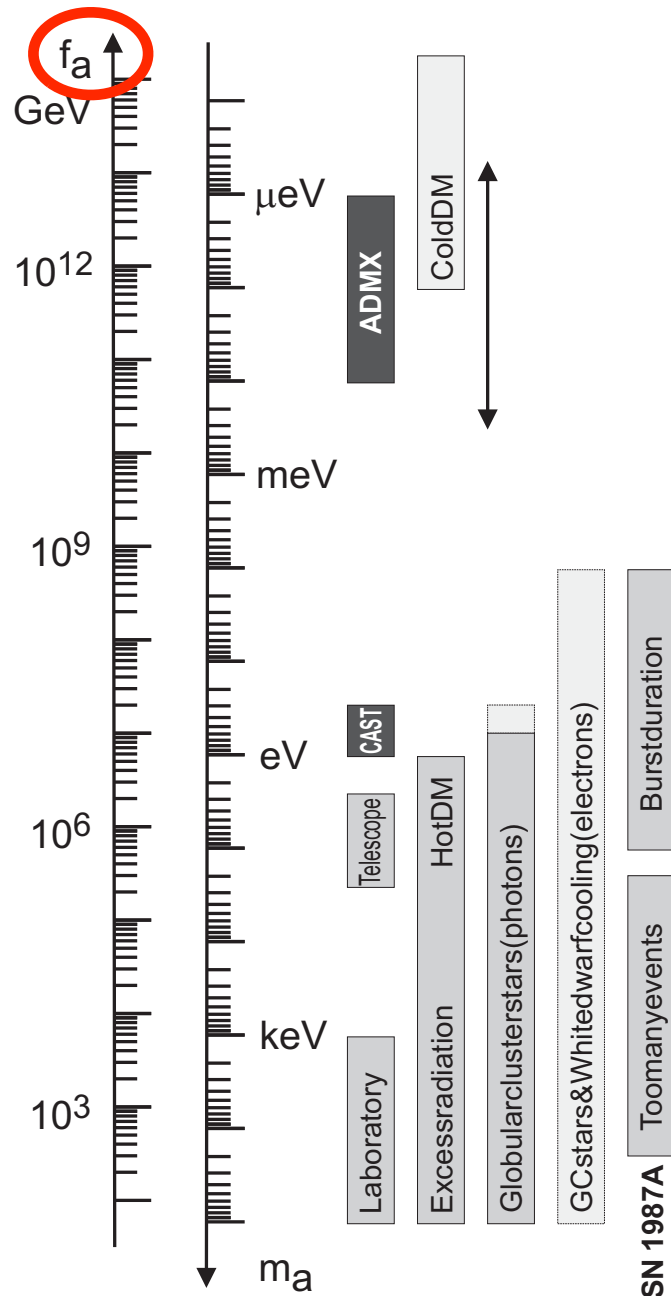
Long-Lived Stau NLSP
[from P. Zalewski's Talk, SUSY 2007]



Probing f_a @ Colliders

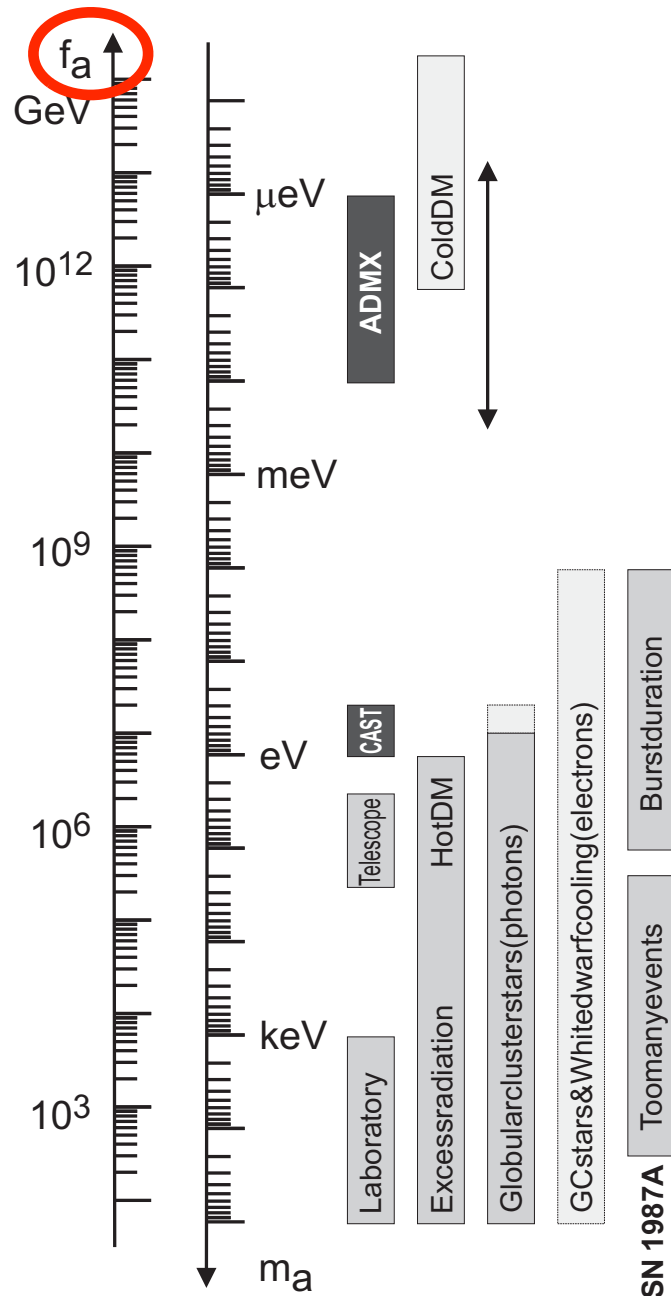
[Brandenburg et al., '05]

Bounds on the Peccei-Quinn Scale



Is the value of the Peccei-Quinn scale inferred from axino searches consistent with astrophysical axion bounds and results from axion searches?

Bounds on the Peccei-Quinn Scale

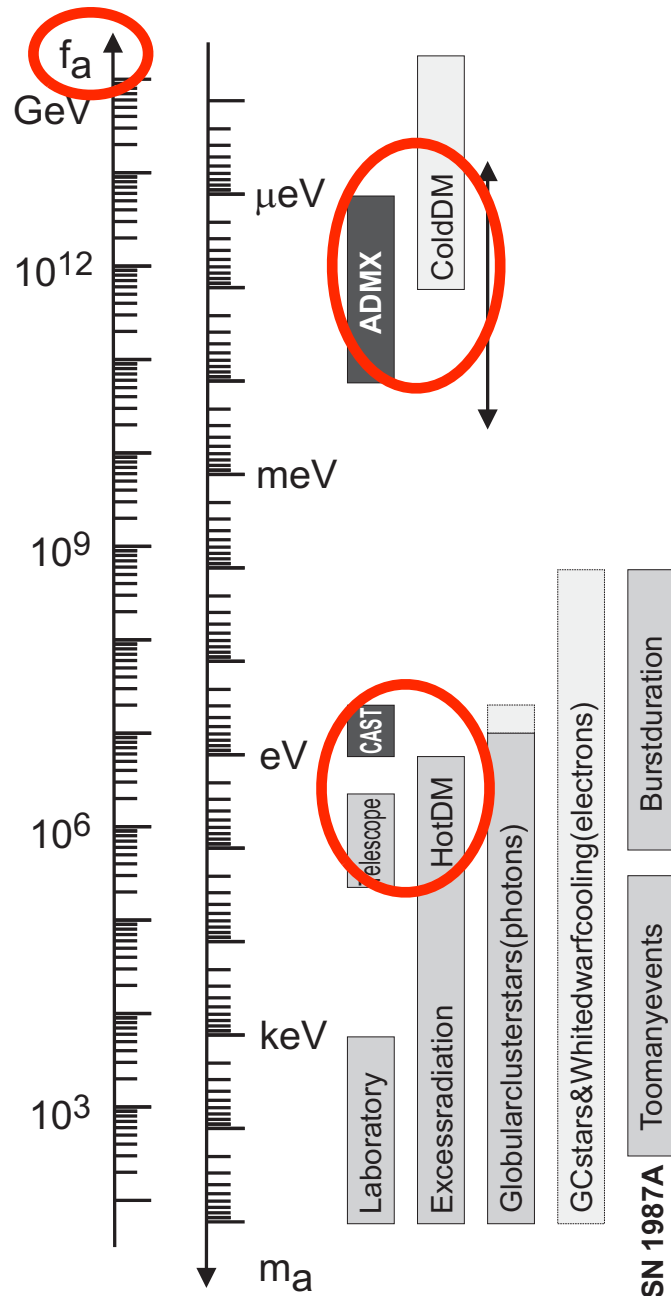


**Agreement between
Axion & Axino Searches**



**Strong Hint for the
Axino LSP**

Bounds on the Peccei-Quinn Scale



**Agreement between
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**Strong Hint for the
Axino LSP**

**Axion DM & Axino DM
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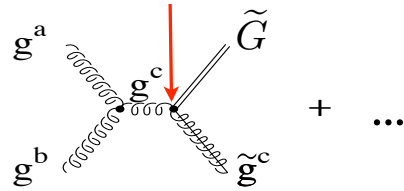
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Thermal Gravitino Production



inflation

radiation dominated

mat. dom.

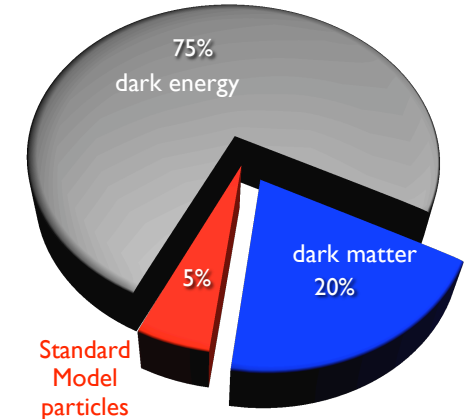
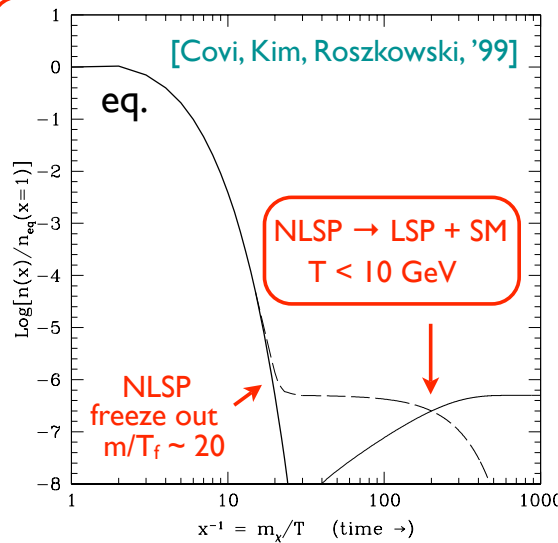
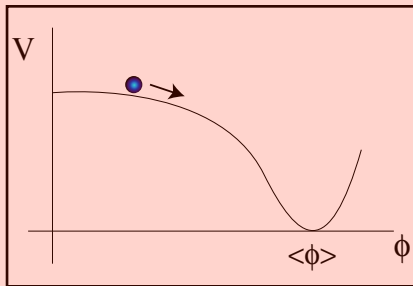
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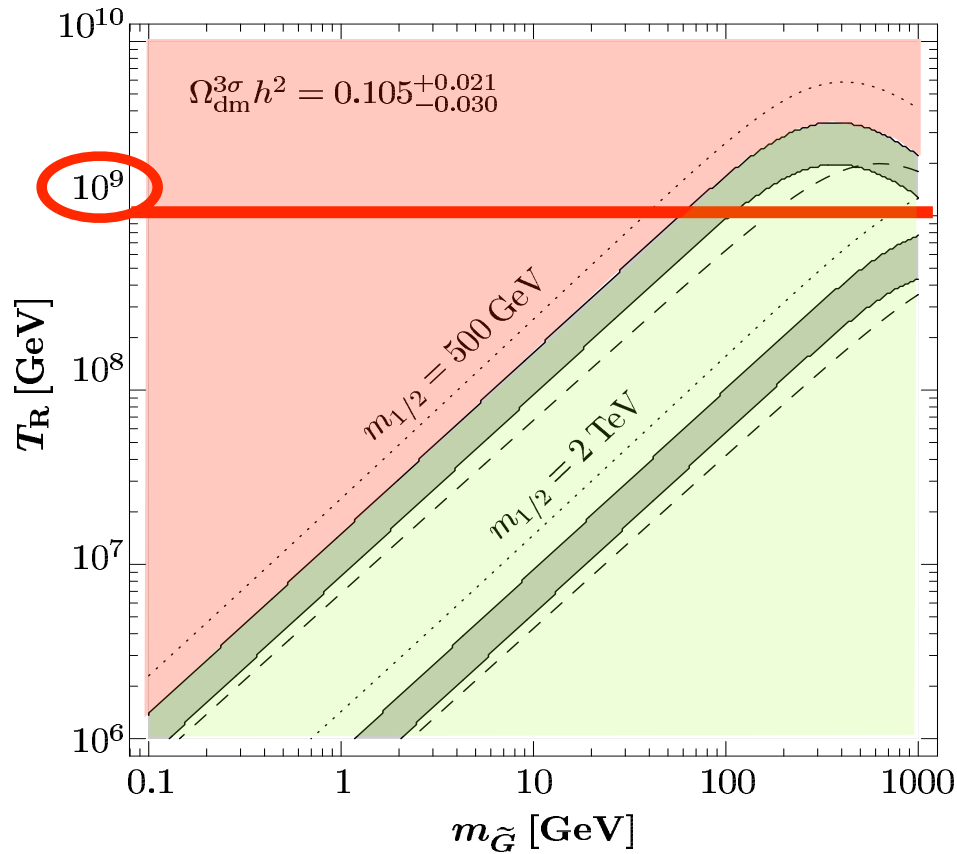
1 eV
CMB

$t_0 = 14 \text{ Gy}$
 $T_0 = 2.73 \text{ K}$
LHC

t
 T
 a

Gravitino LSP Case

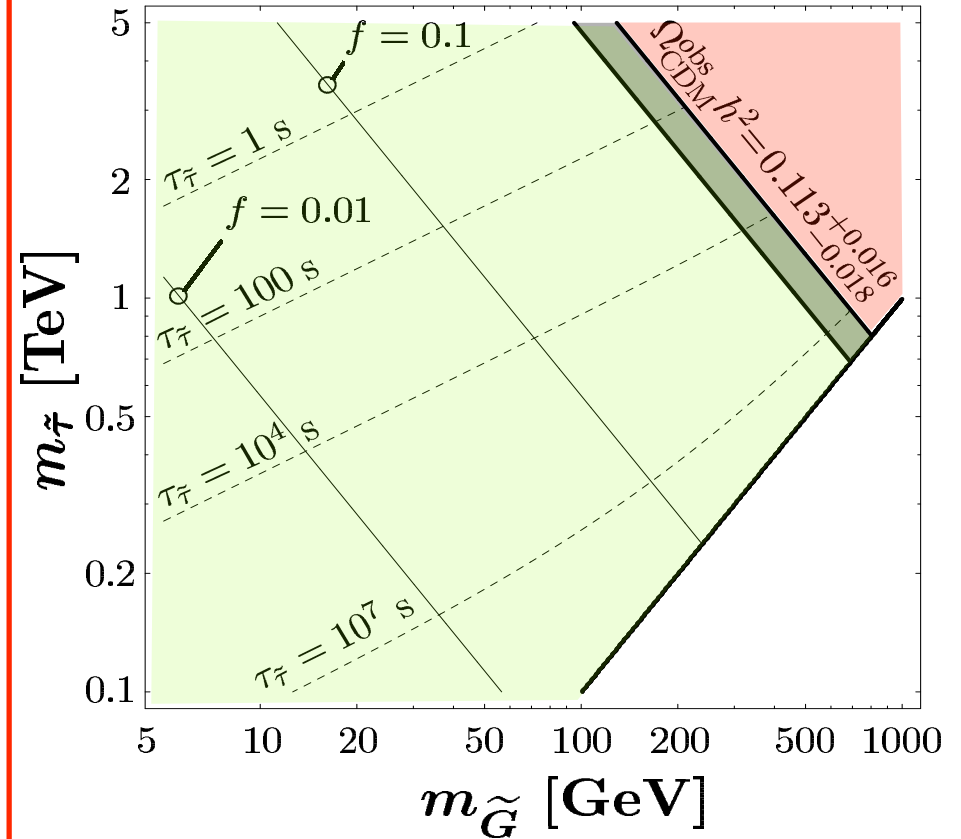
Thermal \tilde{G} Production



[Pradler, FDS, '07]

see also [Moroi, Murayama, Yamaguchi, '93, Asaka, Hamaguchi, Suzuki, '00, Roszkowski et al., '05, Cerdeno et al., '06, FDS '06, Rychkov, Strumia, '07]

$\tilde{\tau}$ NLSP $\rightarrow \tilde{G} + \tau$



[FDS '06]

see also [Borgani, Masiero, Yamaguchi, '96, Asaka, Hamaguchi, Suzuki, '00, Ellis et al., '04, Feng, Su, Takayama, '04]

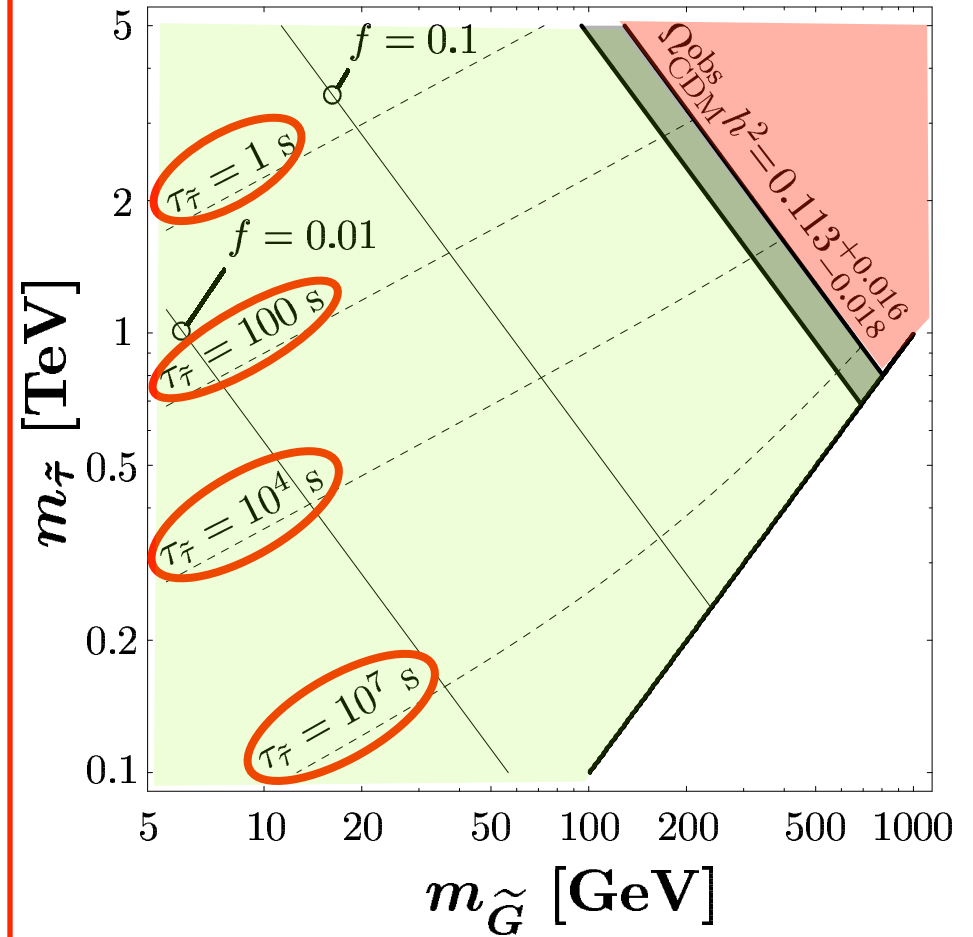
Probing gravitinos experimentally ???



If we are lucky ...

long-lived NLSP

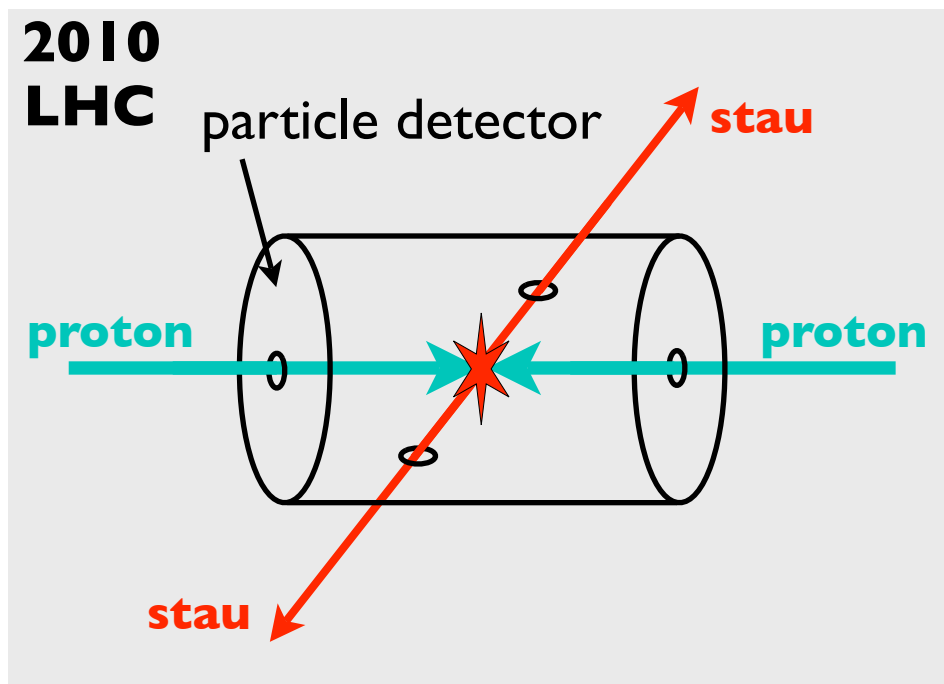
$$\tilde{\tau} \text{ NLSP} \rightarrow \tilde{G} + \tau$$



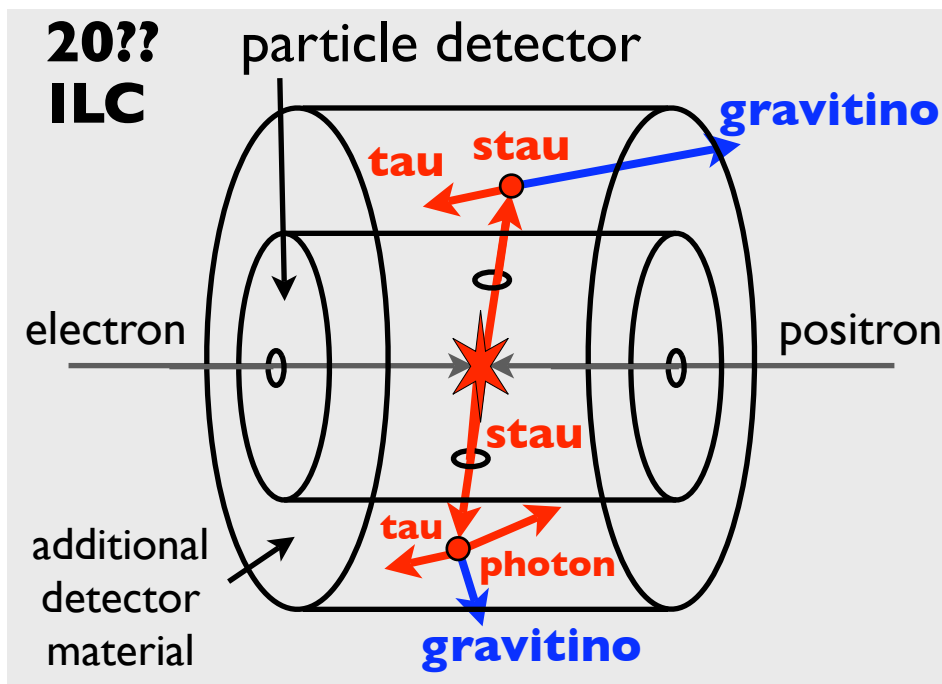
Signatures of Gravitinos in Experiments

- Direct Detection of \tilde{G}
- Direct Production of \tilde{G}

* “stable” charged sparticles



* long-lived charged sparticles



[... ; Buchmüller et al., '04; Hamaguchi et al., '04; Feng, Smith, '05; Martyn, '06; ...]

Summary - Well-motivated DM Candidates

candidate	identity	mass	interactions	production	constraints	experiments
a	axion (spin 0) N.-Goldst. boson PQ symm. break.	$< 0.01 \text{ eV}$	$(p/f_a)^n$ extremely weak $f_a \gtrsim 6 \times 10^8 \text{ GeV}$	misalign. mech.	\leftarrow cold CMB	direct searches with microwave cavities $\hookrightarrow m_a, f_a, g_{a\gamma\gamma}$
$\tilde{\chi}_1^0$ LSP	lightest neutralino (spin 1/2) mixture of $\tilde{B}, \tilde{W}, \tilde{H}_u^0, \tilde{H}_d^0$	$\mathcal{O}(100 \text{ GeV})$	g, g', y_i weak $M_W \sim 100 \text{ GeV}$	therm. relic \tilde{G} decay	\leftarrow cold \leftarrow warm/hot BBN	indirect searches direct searches collider searches $\hookrightarrow m_{\tilde{\chi}_1^0}, \tilde{\chi}_1^0$ coupl.
\tilde{G} LSP	gravitino (spin 3/2) superpartner of the graviton	eV–TeV	$(p/M_P)^n$ extremely weak $M_P = 2.4 \times 10^{18} \text{ GeV}$	therm. prod. NLSP decay	\leftarrow cold \leftarrow warm BBN	$\tilde{\tau}_1$ prod. at colliders + $\tilde{\tau}_1$ collection + $\tilde{\tau}_1$ decay analysis $\hookrightarrow m_{\tilde{G}}, M_P (?), T_R$
\tilde{a} LSP	axino (spin 1/2) superpartner of the axion	eV–GeV	$(p/f_a)^n$ extremely weak $f_a \gtrsim 6 \times 10^8 \text{ GeV}$	therm. prod. NLSP decay	\leftarrow cold/warm \leftarrow warm/hot BBN	$\tilde{\tau}_1$ prod. at colliders + $\tilde{\tau}_1$ collection + $\tilde{\tau}_1$ decay analysis $\hookrightarrow m_{\tilde{a}} (?), f_a, T_R (?)$

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a

$\tilde{\chi}_1^0$ LSP

\tilde{G} LSP

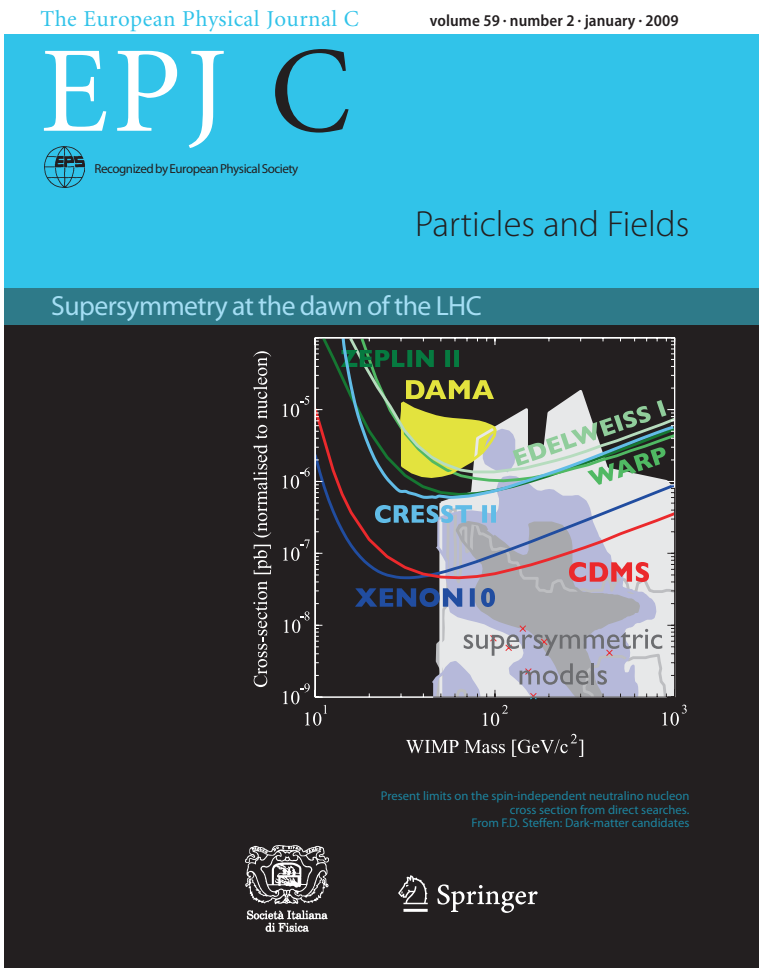
\tilde{a} LSP

For a review (including an extensive list of references),

see

[FDS, Dark Matter Candidates, Eur. Phys. J. C59 (2009) 557, arXiv:0811.3347]

in



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Scenario I - Axion CDM (+ SUSY DM)

candidate	identity	mass	interactions	production	constraints	experiments events
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\tilde{G} LSP	gravitino (spin 3/2) superpartner of the graviton	eV–TeV	$(p/M_P)^n$ extremely weak $M_P = 2.4 \times 10^{18} \text{ GeV}$	therm. prod. NLSP decay	\leftarrow cold \leftarrow warm BBN	$\tilde{\tau}_1$ prod. at colliders + $\tilde{\tau}_1$ collection + $\tilde{\tau}_1$ decay analysis $\hookrightarrow m_{\tilde{G}}, M_P (?), T_R$
\tilde{a} LSP	axino (spin 1/2) superpartner of the axion	eV–GeV	$(p/f_a)^n$ extremely weak $f_a \gtrsim 6 \times 10^8 \text{ GeV}$	therm. prod. NLSP decay	\leftarrow cold/warm \leftarrow warm/hot BBN	$\tilde{\tau}_1$ prod. at colliders + $\tilde{\tau}_1$ collection + $\tilde{\tau}_1$ decay analysis $\hookrightarrow m_{\tilde{a}} (?), f_a, T_R (?)$

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still viable

Scenario 2 - WIMP DM (+ Axion DM)

candidate	identity	mass	interactions	production	constraints	experiments
a	axion (spin 0) N.-Goldst. boson PQ symm. break.	$< 0.01 \text{ eV}$	$(p/f_a)^n$ extremely weak $f_a \gtrsim 6 \times 10^8 \text{ GeV}$	misalign. mech.	\leftarrow cold CMB	direct searches with microwave cavities $\hookrightarrow m_a, f_a, g_{a\gamma\gamma}$
$\tilde{\chi}_1^0$ LSP	lightest neutralino (spin 1/2) mixture of $\tilde{B}, \tilde{W}, \tilde{H}_u^0, \tilde{H}_d^0$	$\mathcal{O}(100 \text{ GeV})$	g, g', y_i weak $M_W \sim 100 \text{ GeV}$	therm. relic \tilde{G} decay	\leftarrow cold \leftarrow warm/hot BBN	indirect searches direct searches events collider searches $\hookrightarrow m_{\tilde{\chi}_1^0}, \tilde{\chi}_1^0$ coupl.
\tilde{G} LSP	gravitino (spin 3/2) superpartner of the graviton	eV–TeV	$(p/M_P)^n$ extremely weak $M_P = 2.4 \times 10^{18} \text{ GeV}$	therm. prod. NLSP decay	\leftarrow cold \leftarrow warm BBN	$\tilde{\tau}_1$ prod. at colliders + $\tilde{\tau}_1$ collection + $\tilde{\tau}_1$ decay analysis $\hookrightarrow m_{\tilde{G}}, M_P (?), T_R$
\tilde{a} LSP	axino (spin 1/2) superpartner of the axion	eV–GeV	$(p/f_a)^n$ extremely weak $f_a \gtrsim 6 \times 10^8 \text{ GeV}$	therm. prod. NLSP decay	\leftarrow cold/warm \leftarrow warm/hot BBN	$\tilde{\tau}_1$ prod. at colliders + $\tilde{\tau}_1$ collection + $\tilde{\tau}_1$ decay analysis $\hookrightarrow m_{\tilde{a}} (?), f_a, T_R (?)$

Scenario 2 - WIMP DM (+ Axion DM)

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$\tilde{\chi}_1^0$ LSP	lightest neutralino (spin 1/2) mixture of $\tilde{B}, \tilde{W}, \tilde{H}_u^0, \tilde{H}_d^0$	$\mathcal{O}(100 \text{ GeV})$	g, g', y_i weak $M_W \sim 100 \text{ GeV}$	therm. relic \tilde{G} decay	\leftarrow cold \leftarrow warm/hot BBN	indirect searches direct searches events collider searches $\hookrightarrow m_{\tilde{\chi}_1^0}, \tilde{\chi}_1^0$ coupl.
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Scenario 3 - EWIP DM (+ Axion DM)

candidate	identity	mass	interactions	production	constraints	experiments
a	axion (spin 0) N.-Goldst. boson PQ symm. break.	$< 0.01 \text{ eV}$	$(p/f_a)^n$ extremely weak $f_a \gtrsim 6 \times 10^8 \text{ GeV}$	misalign. mech.	\leftarrow cold CMB	direct searches with microwave cavities $\hookrightarrow m_a, f_a, g_{a\gamma\gamma}$
$\tilde{\chi}_1^0$ LSP	lightest neutralino (spin 1/2) mixture of $\tilde{B}, \tilde{W}, \tilde{H}_u^0, \tilde{H}_d^0$	$\mathcal{O}(100 \text{ GeV})$	g, g', y_i weak $M_W \sim 100 \text{ GeV}$	therm. relic \tilde{G} decay	\leftarrow cold \leftarrow warm/hot BBN	indirect searches direct searches collider searches $\hookrightarrow m_{\tilde{\chi}_1^0}, \tilde{\chi}_1^0$ coupl. events
\tilde{G} LSP	gravitino (spin 3/2) superpartner of the graviton	eV–TeV	$(p/M_P)^n$ extremely weak $M_P = 2.4 \times 10^{18} \text{ GeV}$	therm. prod. NLSP decay	\leftarrow cold \leftarrow warm BBN	$\tilde{\tau}_1$ prod. at colliders + $\tilde{\tau}_1$ collection + $\tilde{\tau}_1$ decay analysis $\hookrightarrow m_{\tilde{G}}, M_P (?), T_R$
\tilde{a} LSP	axino (spin 1/2) superpartner of the axion	eV–GeV	$(p/f_a)^n$ extremely weak $f_a \gtrsim 6 \times 10^8 \text{ GeV}$	therm. prod. NLSP decay	\leftarrow cold/warm \leftarrow warm/hot BBN	$\tilde{\tau}_1$ prod. at colliders + $\tilde{\tau}_1$ collection events + $\tilde{\tau}_1$ decay analysis $\hookrightarrow m_{\tilde{a}} (?), f_a, T_R (?)$

Scenario 3 - EWIP DM (+ Axion DM)

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Conclusion

To clarify the (particle ?) identity of dark matter, it will be crucial to have experimental & obs. data from the many complementary approaches: direct, indirect & collider dm searches